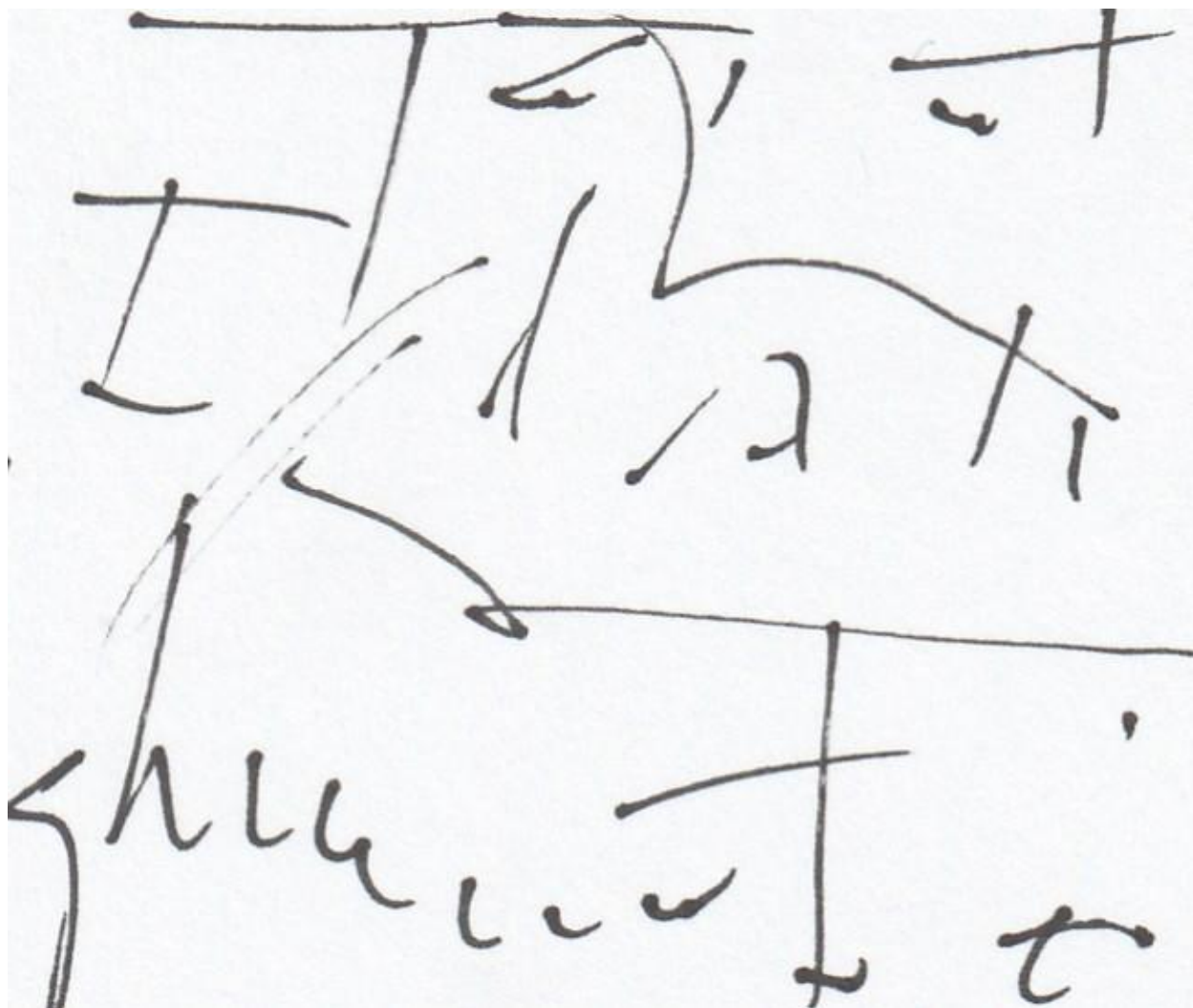
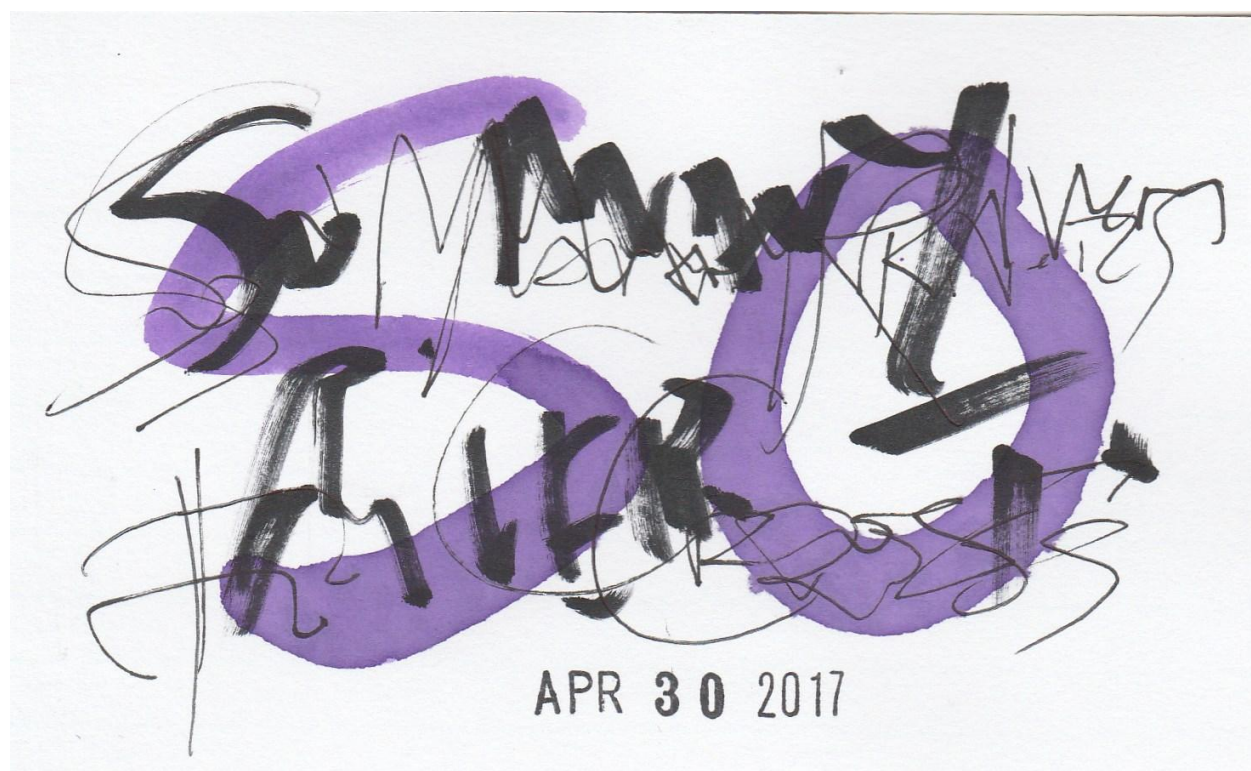
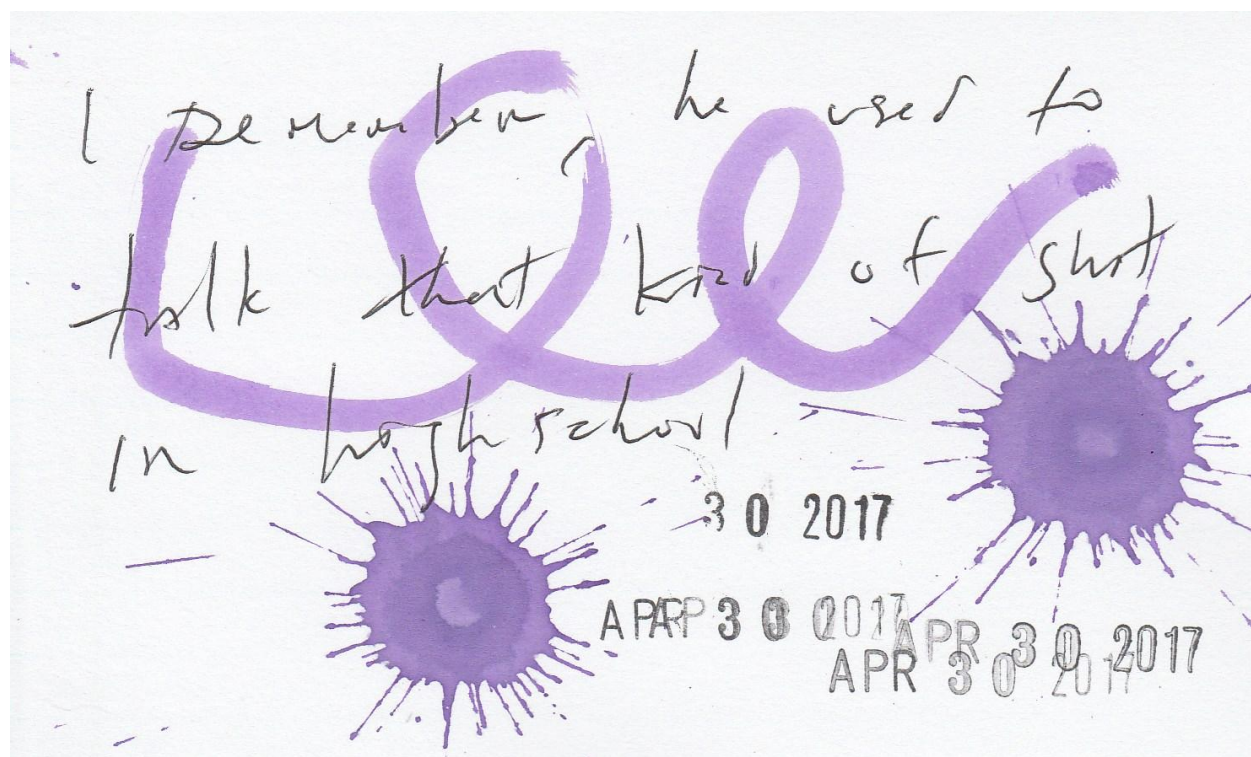


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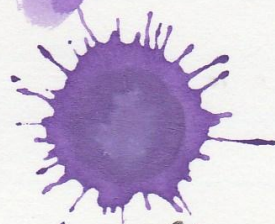
visual poems ongoing research 2017 -vol. 5





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APR 29 2017



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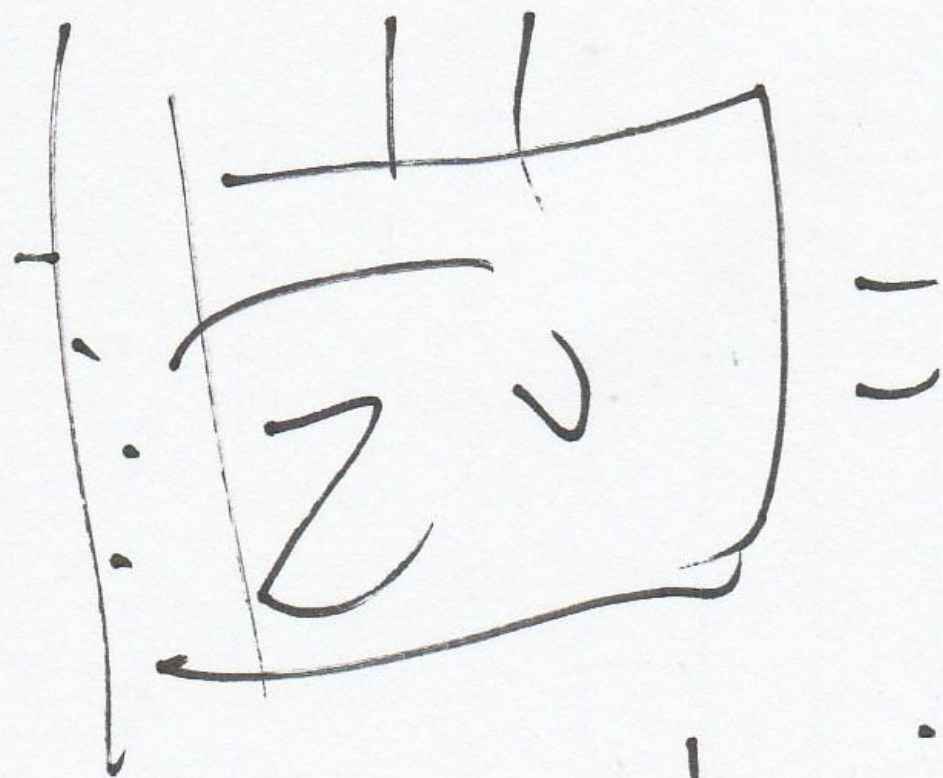
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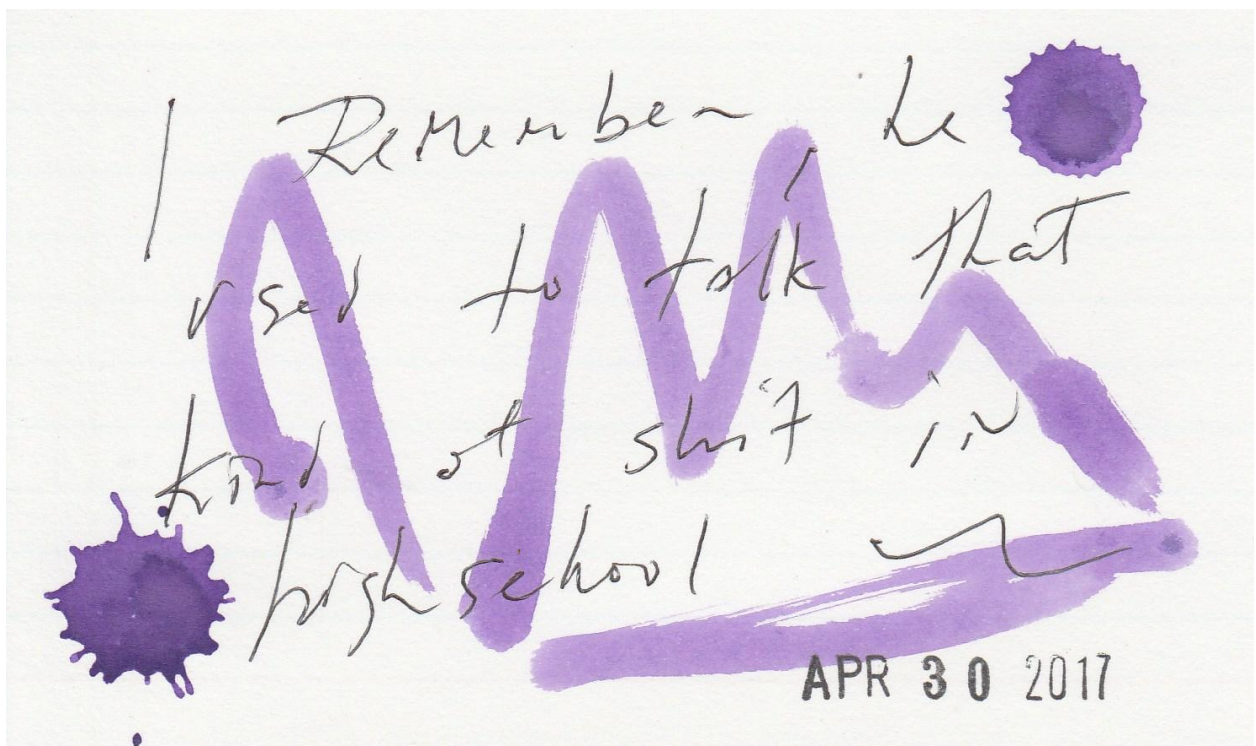




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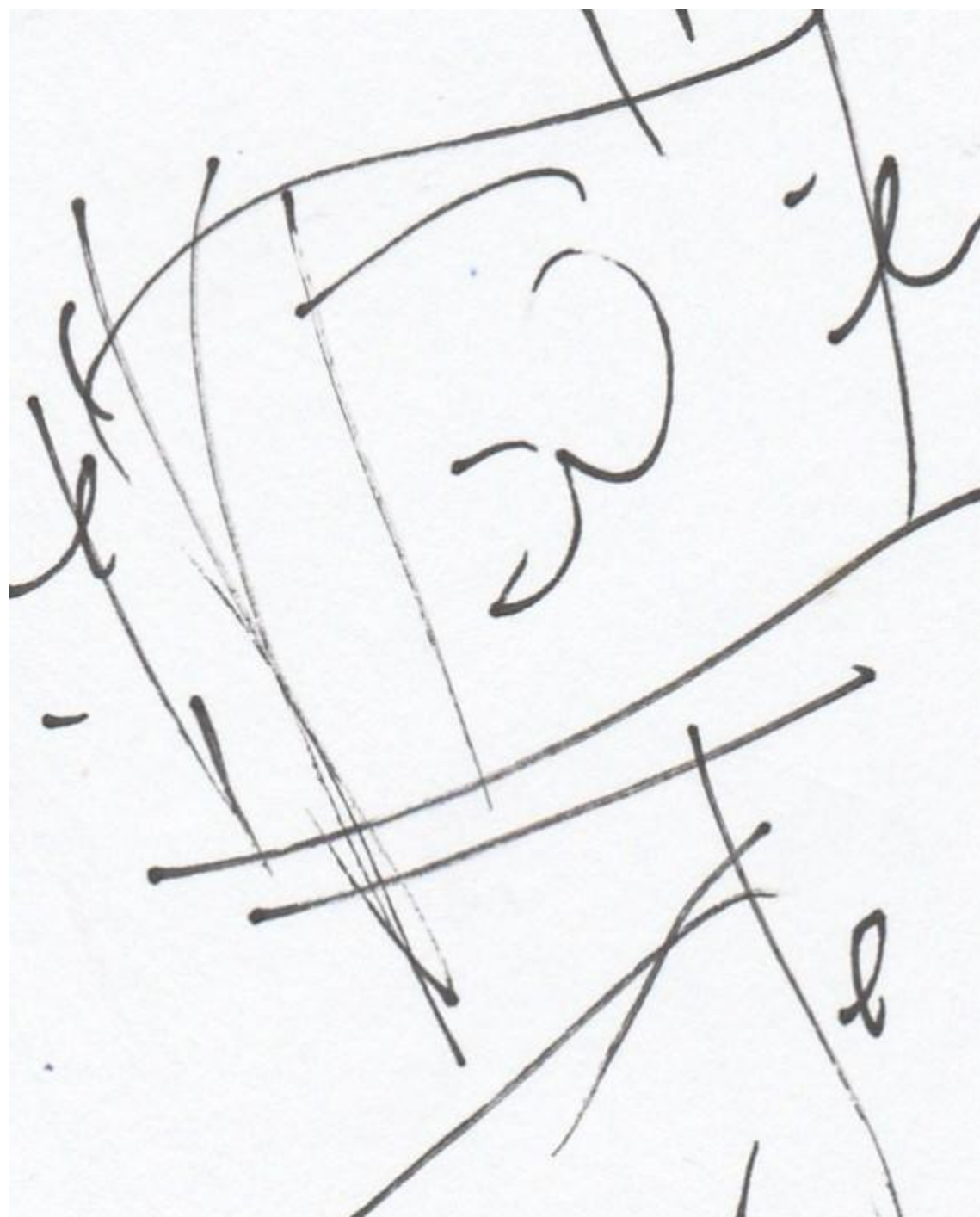




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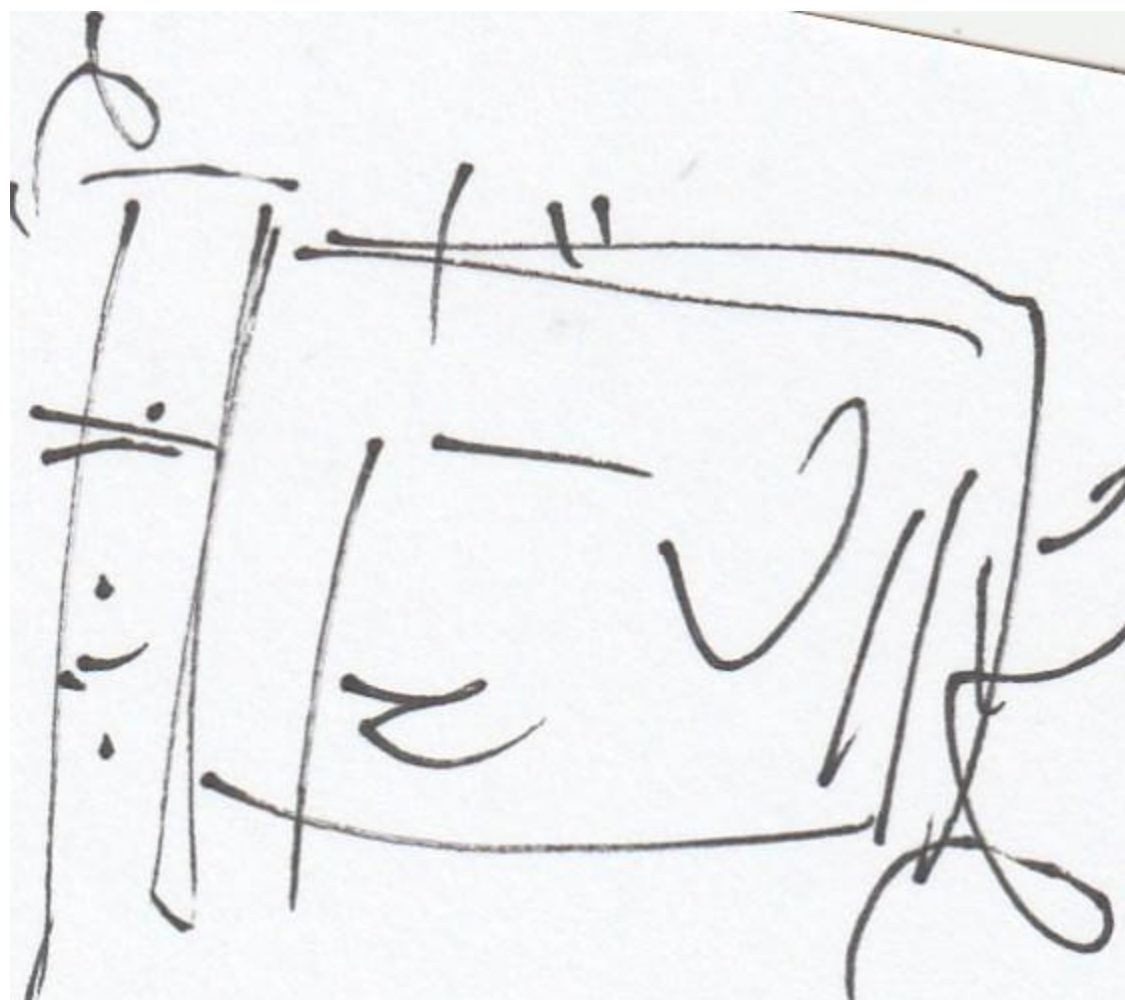
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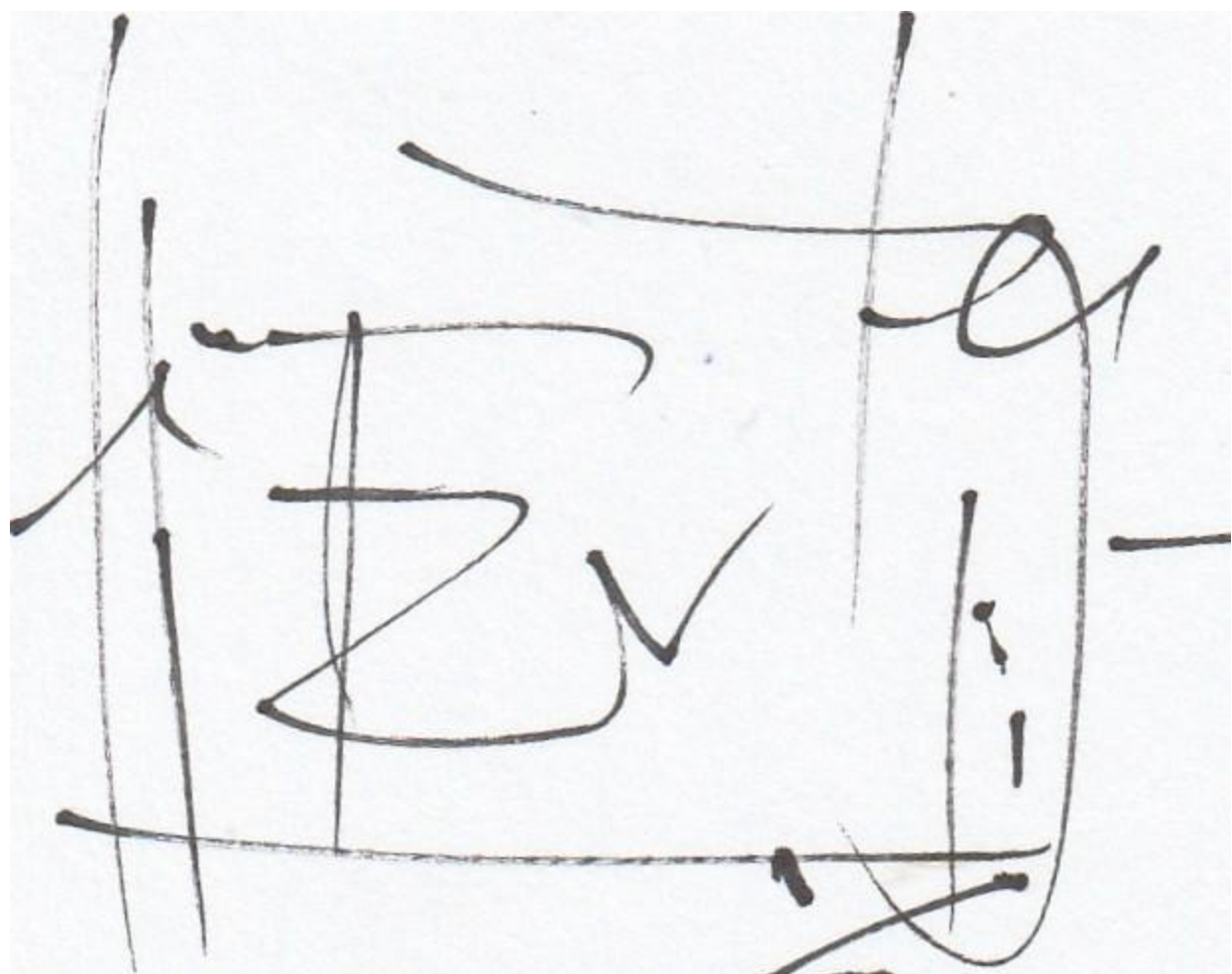


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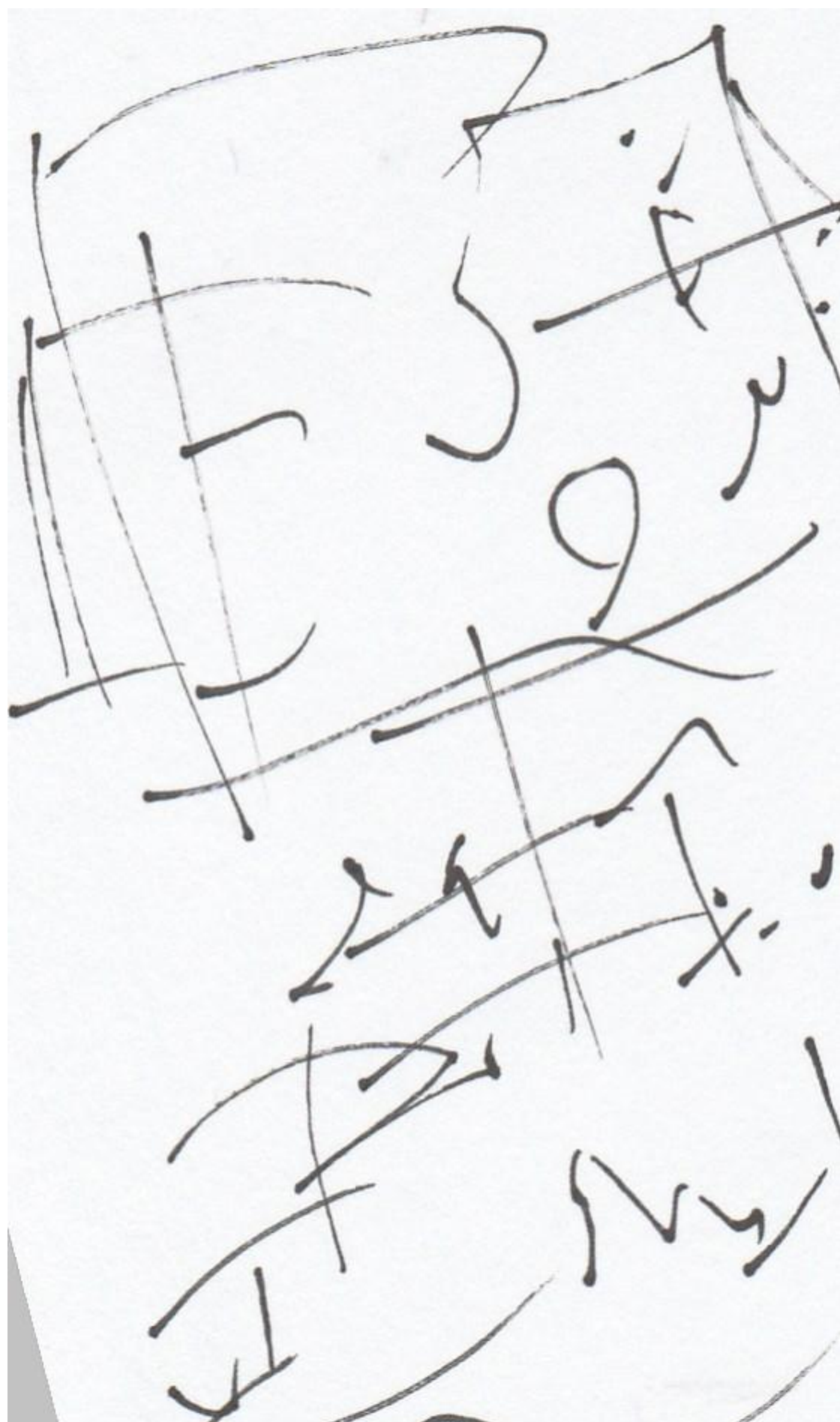


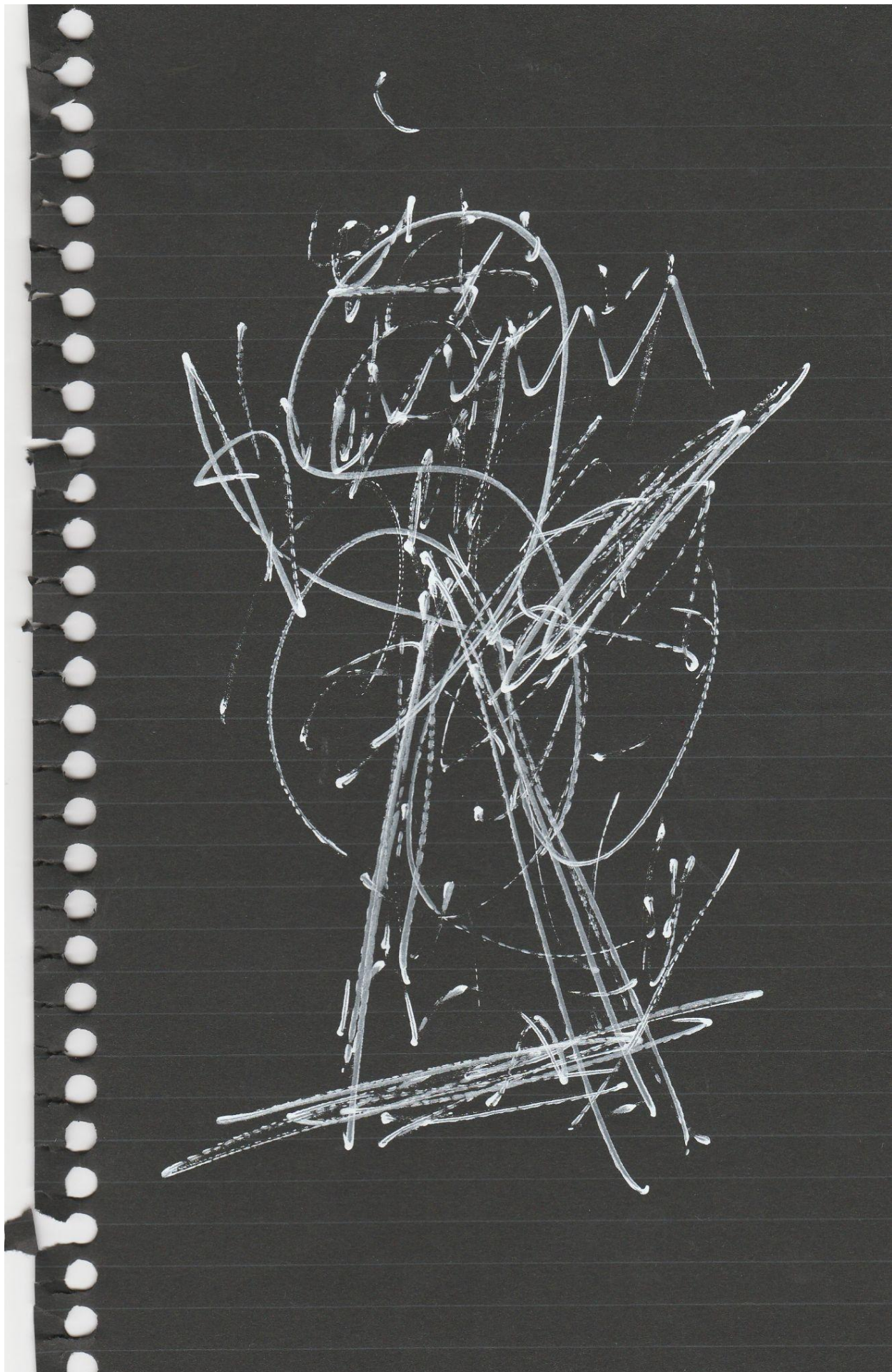




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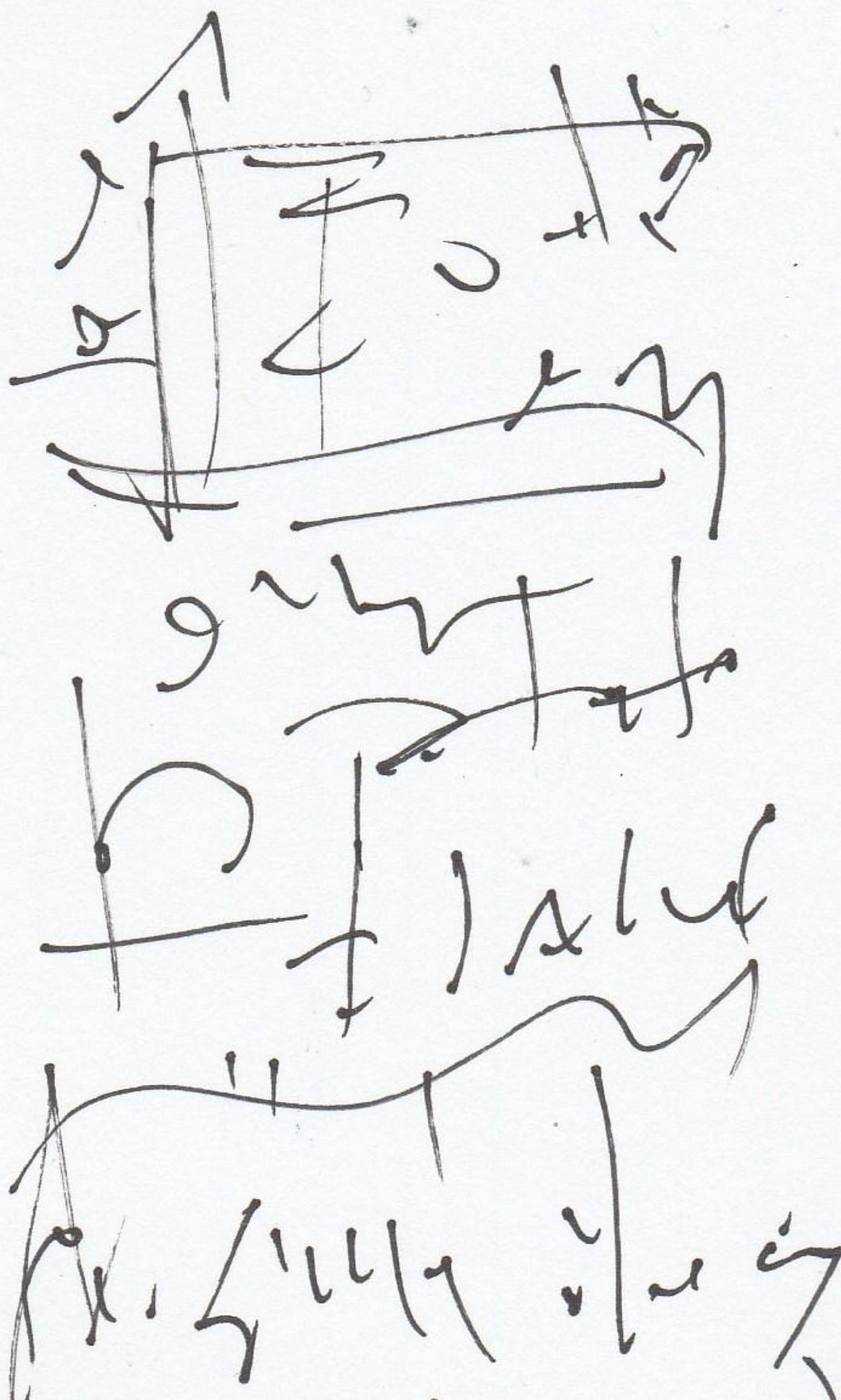








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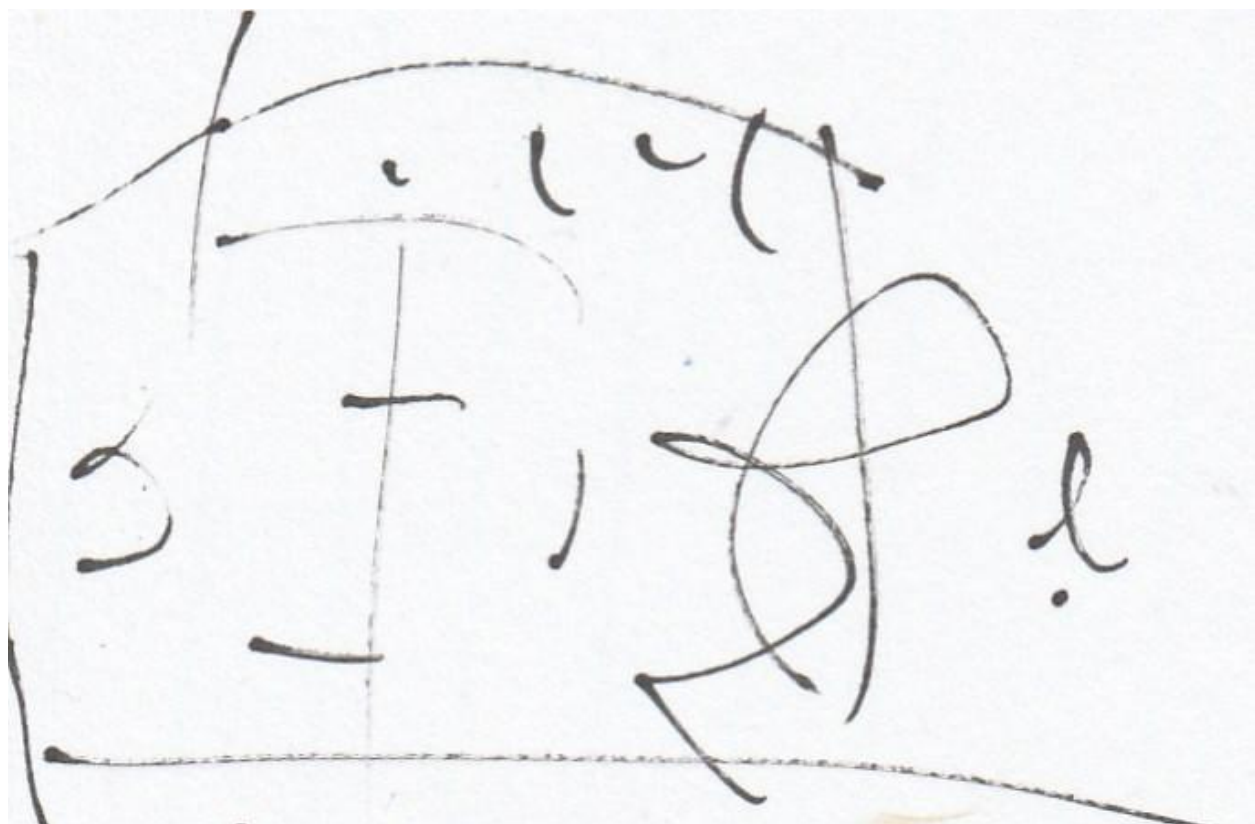
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What does it say?  
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APR 30 2017

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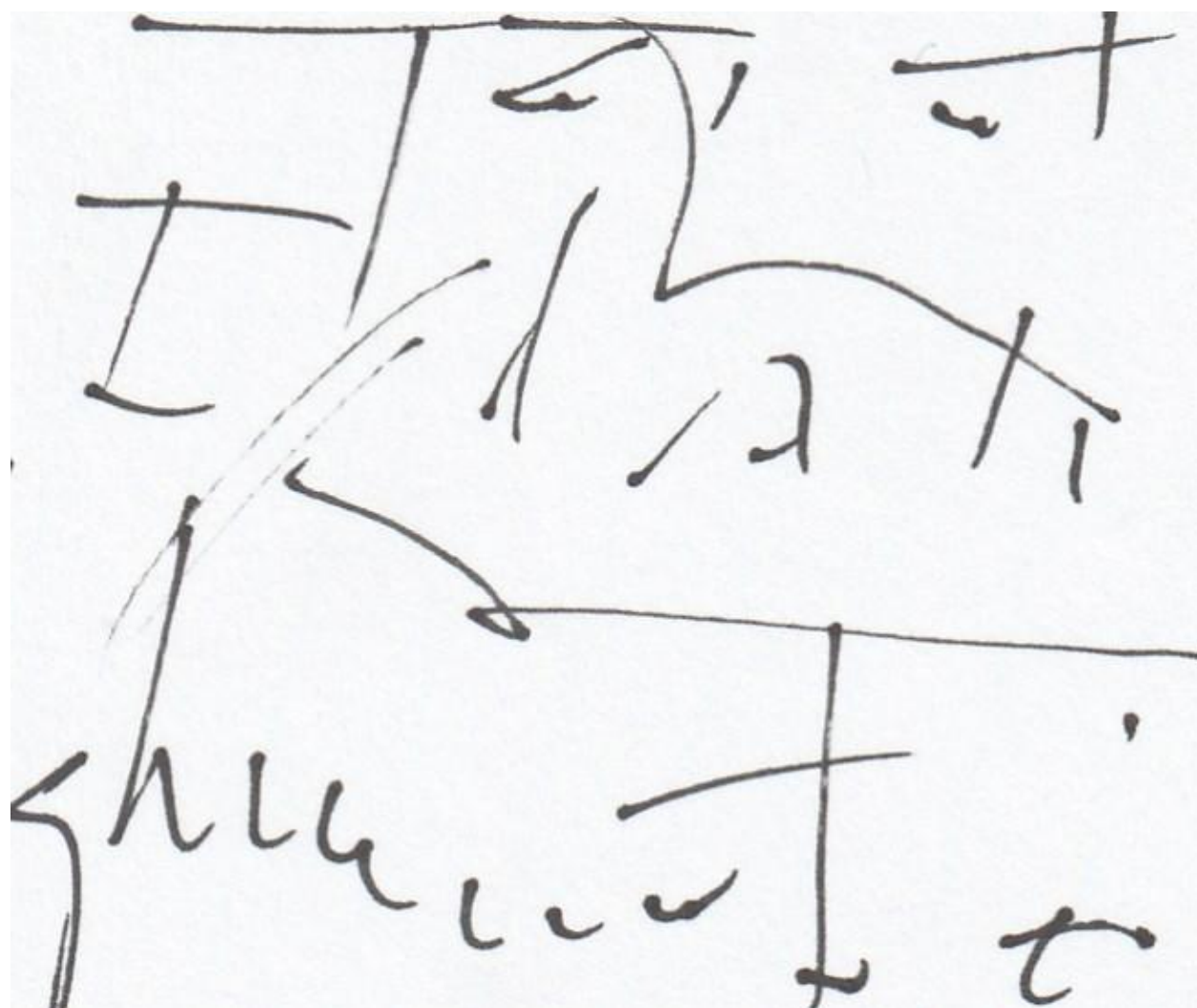




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Handwritten scribble in black ink, possibly resembling the word "LOAN" or a signature.

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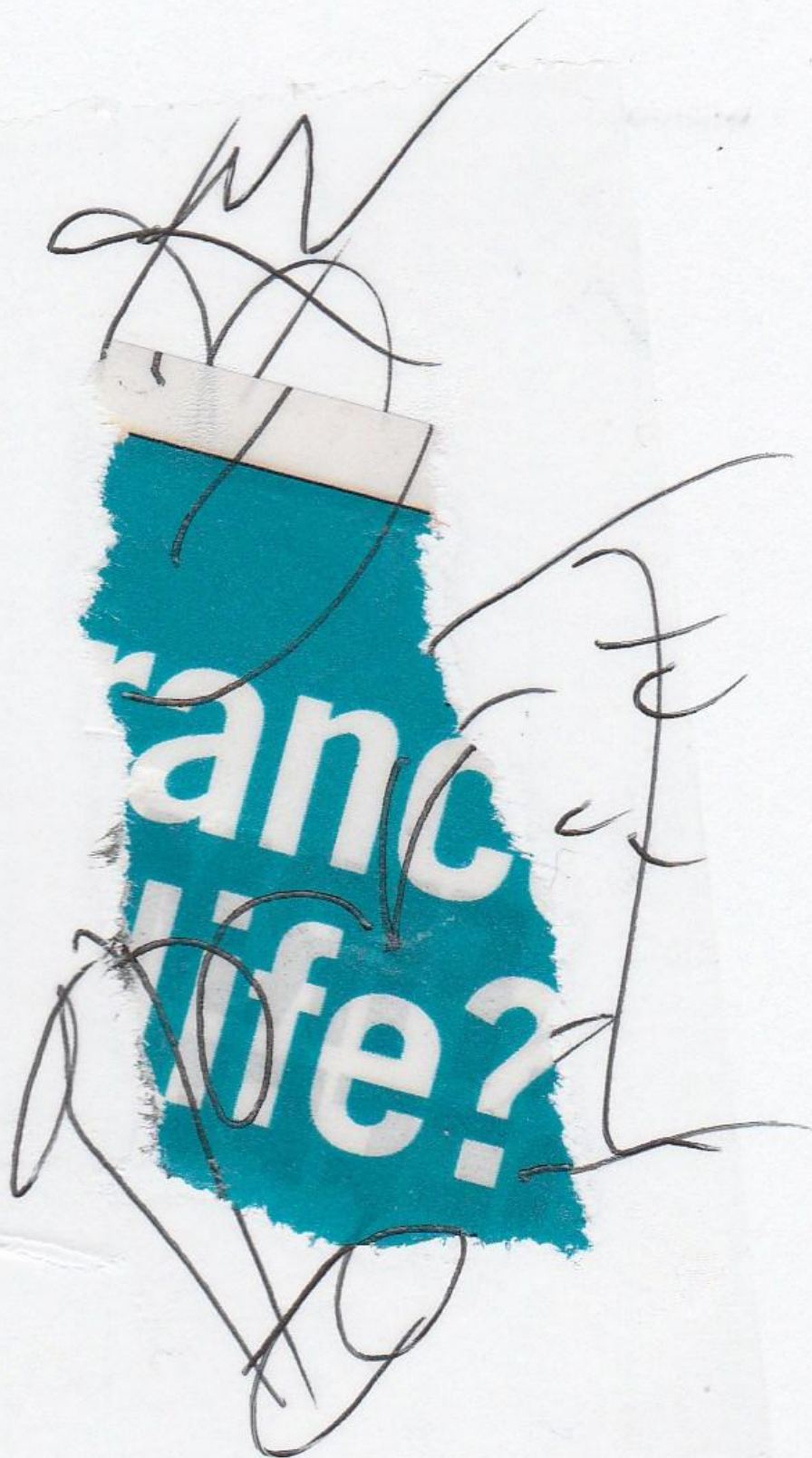
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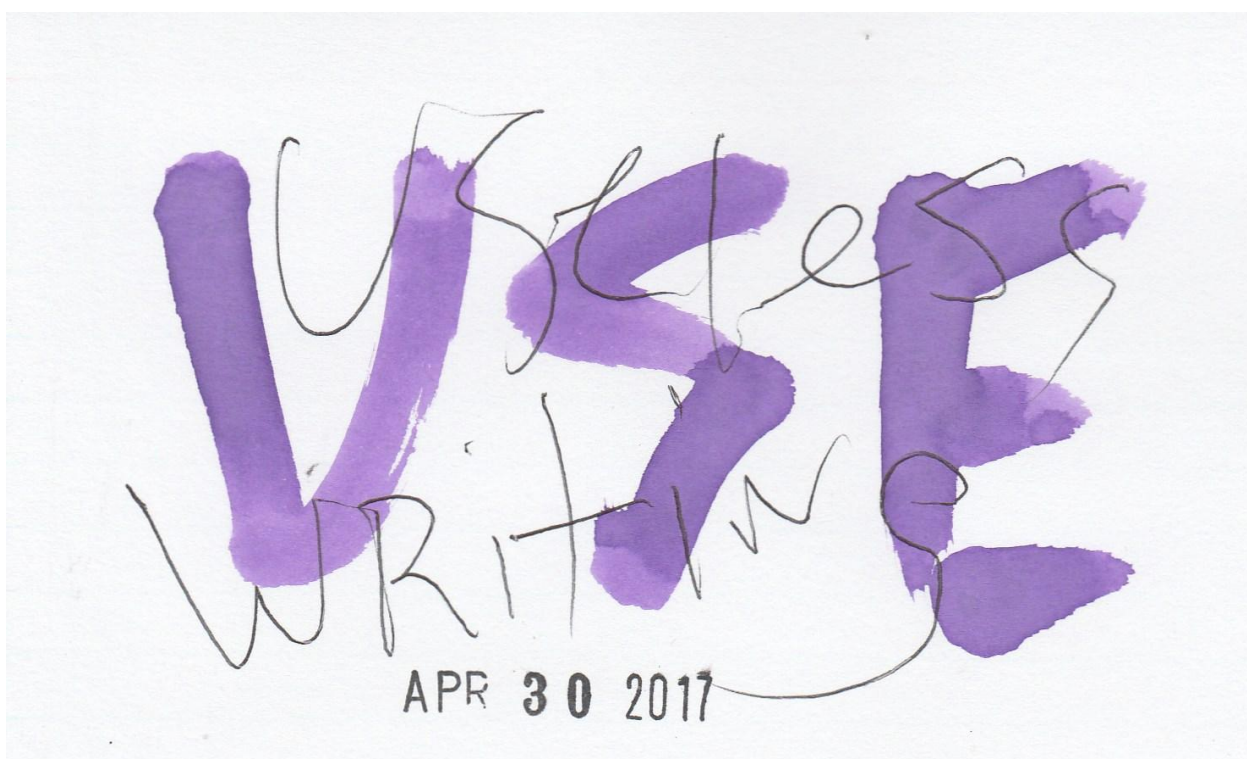
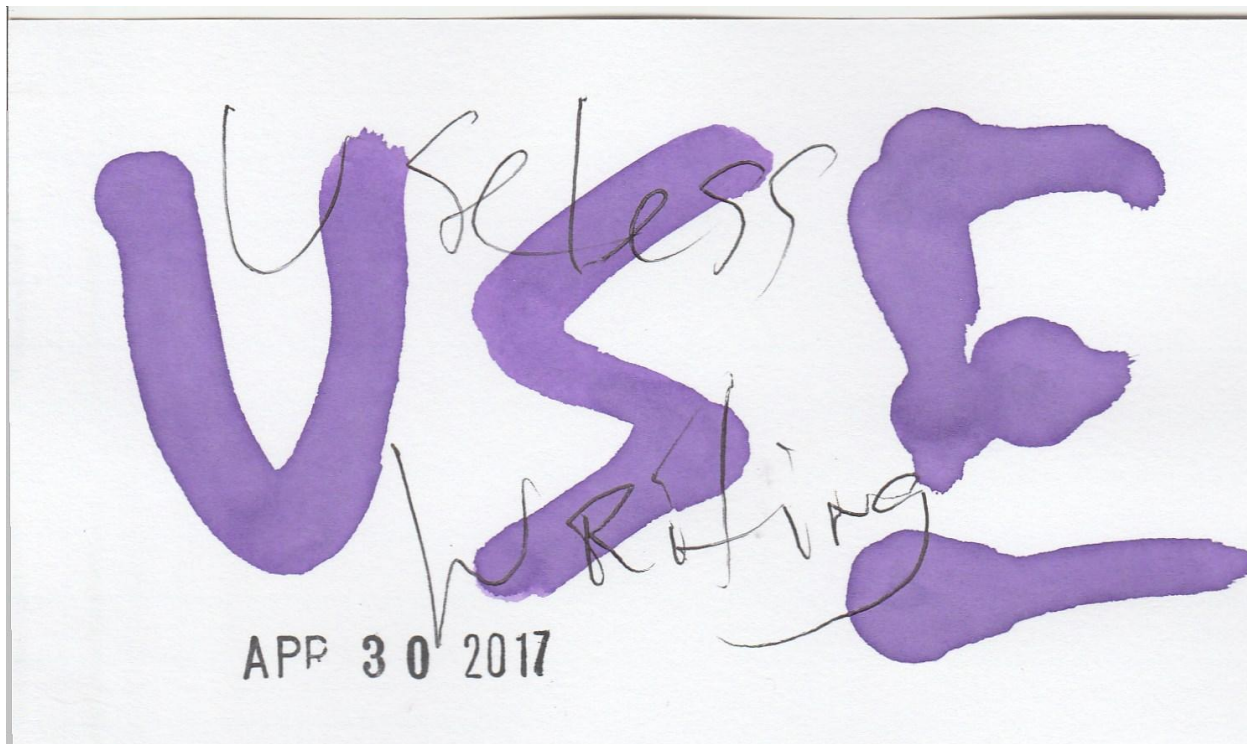


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AP 30 2017



refrigerator telephonic meander refer



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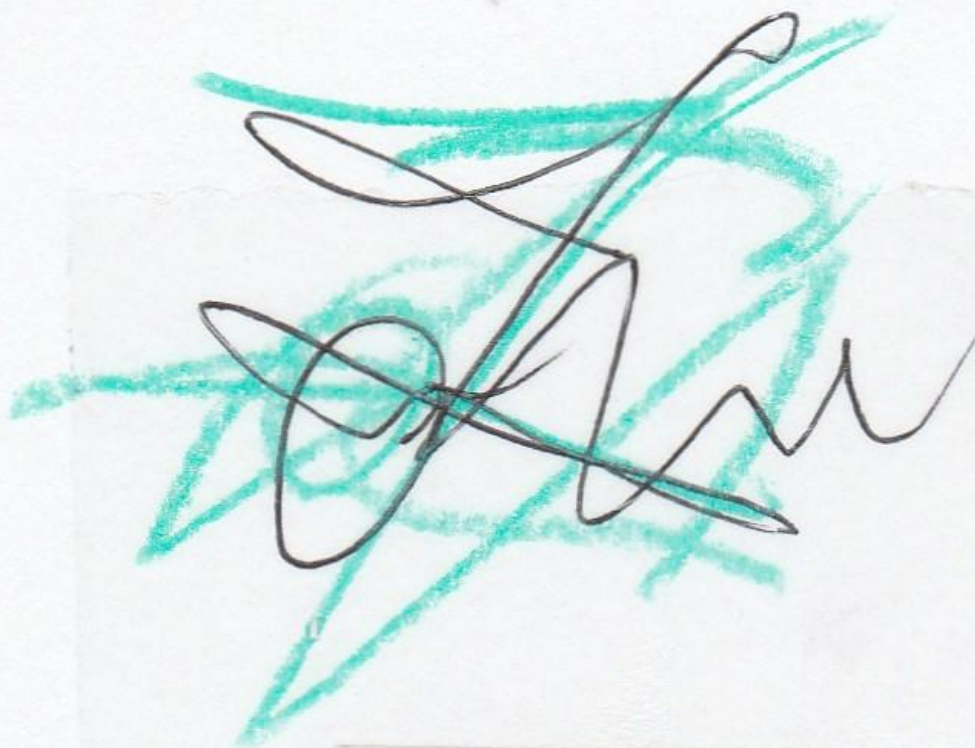
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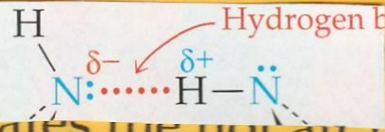
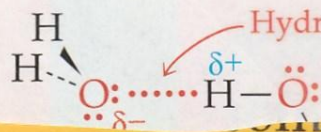
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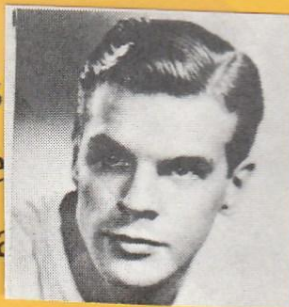


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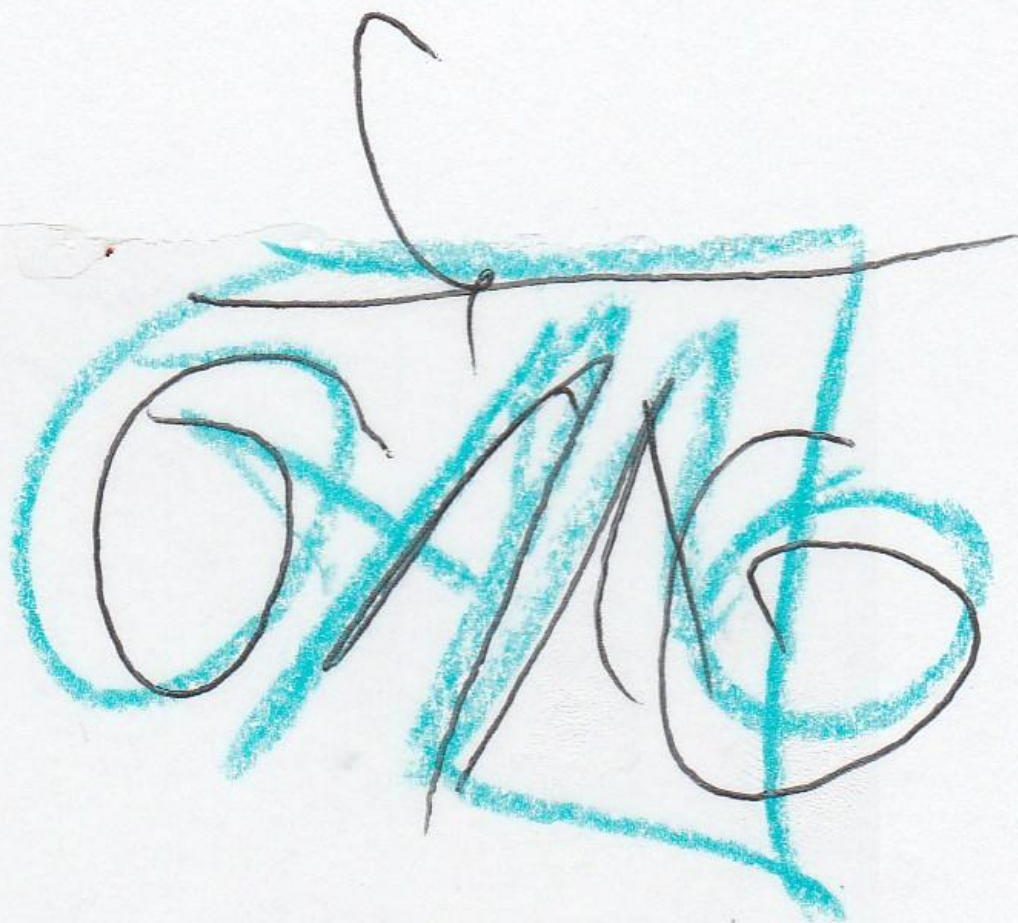
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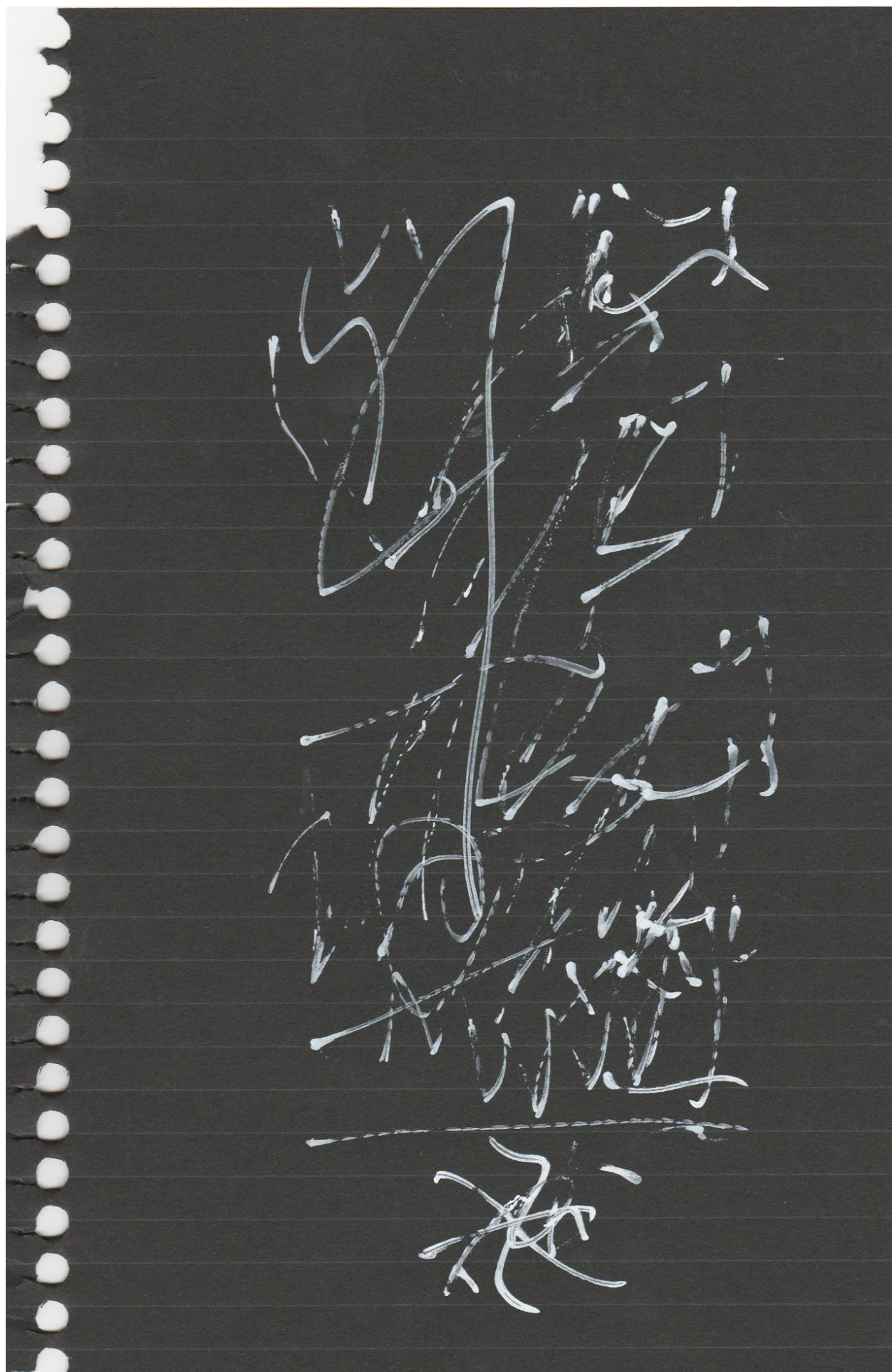


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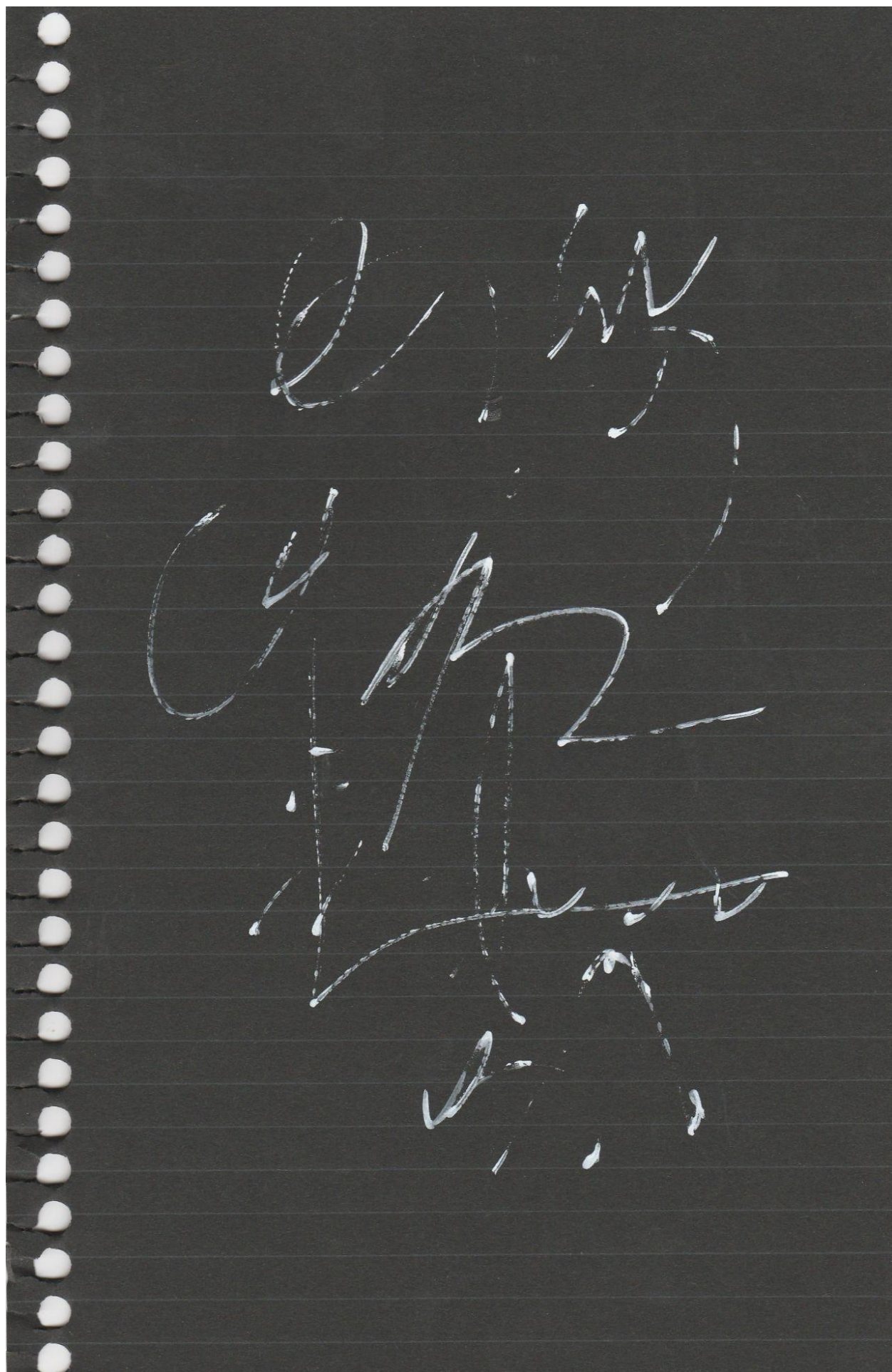
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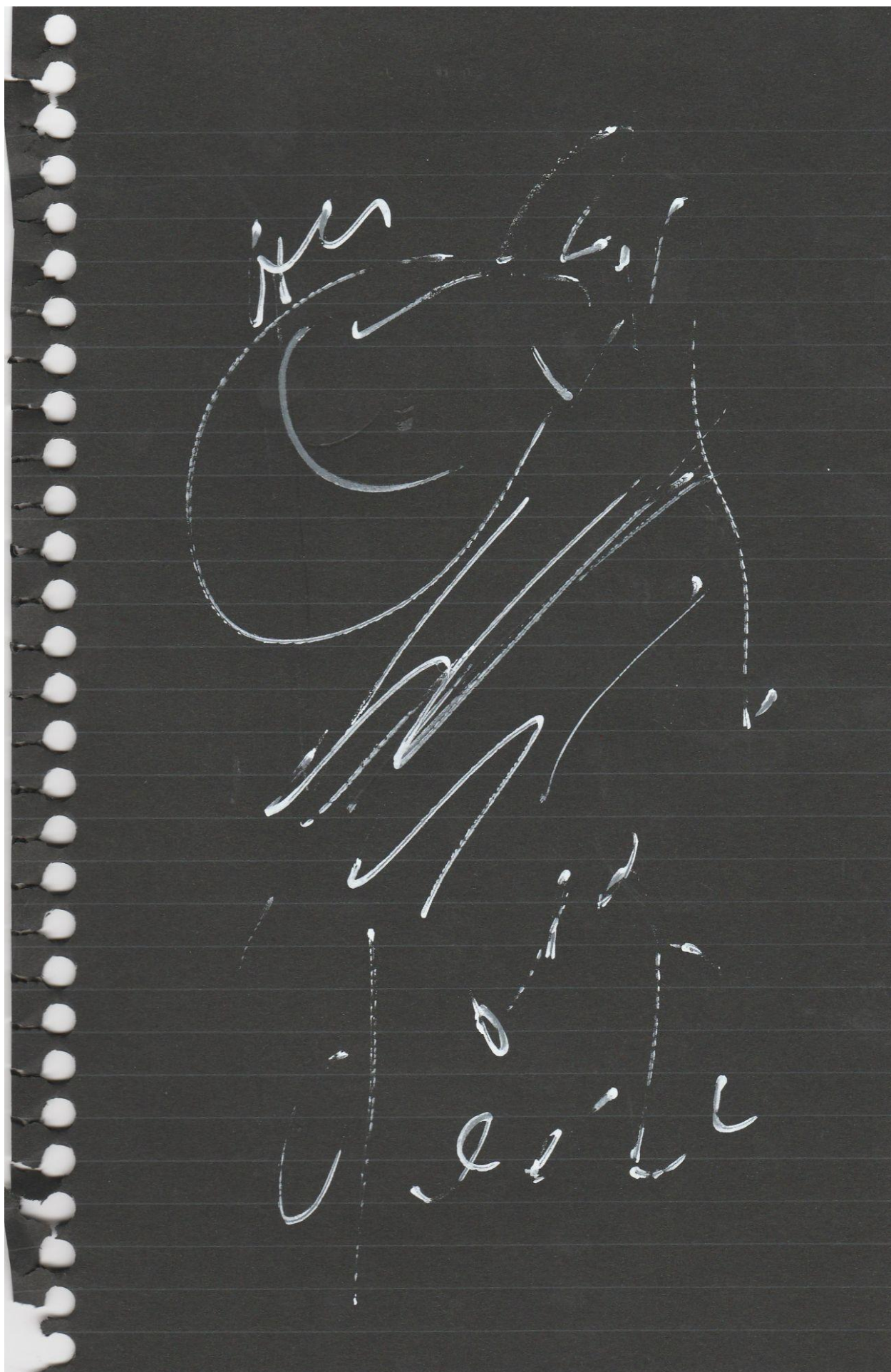




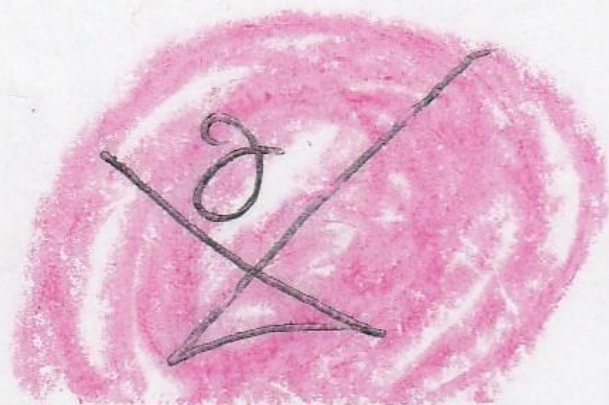






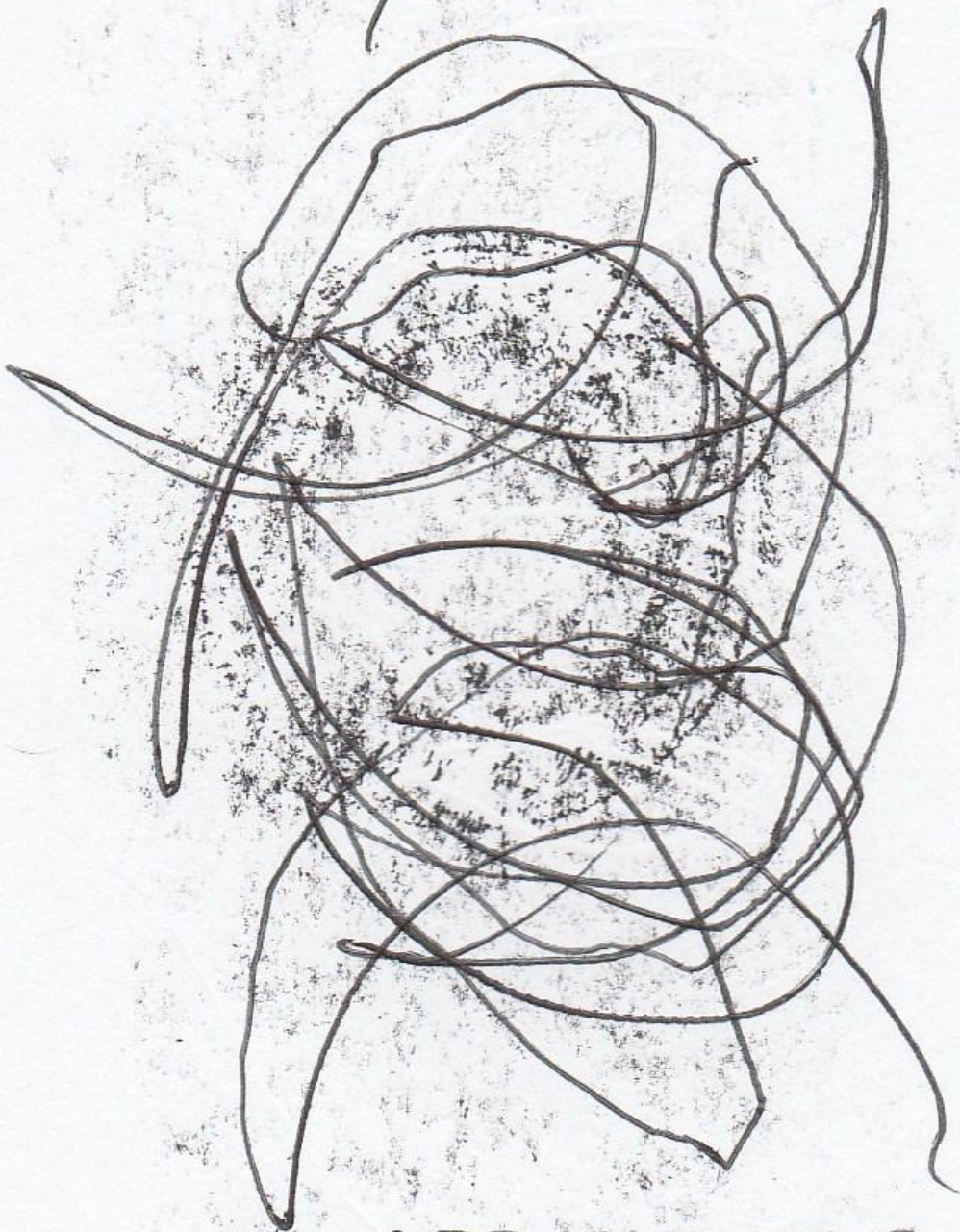






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TRANSACTIONS



APR 30 2017



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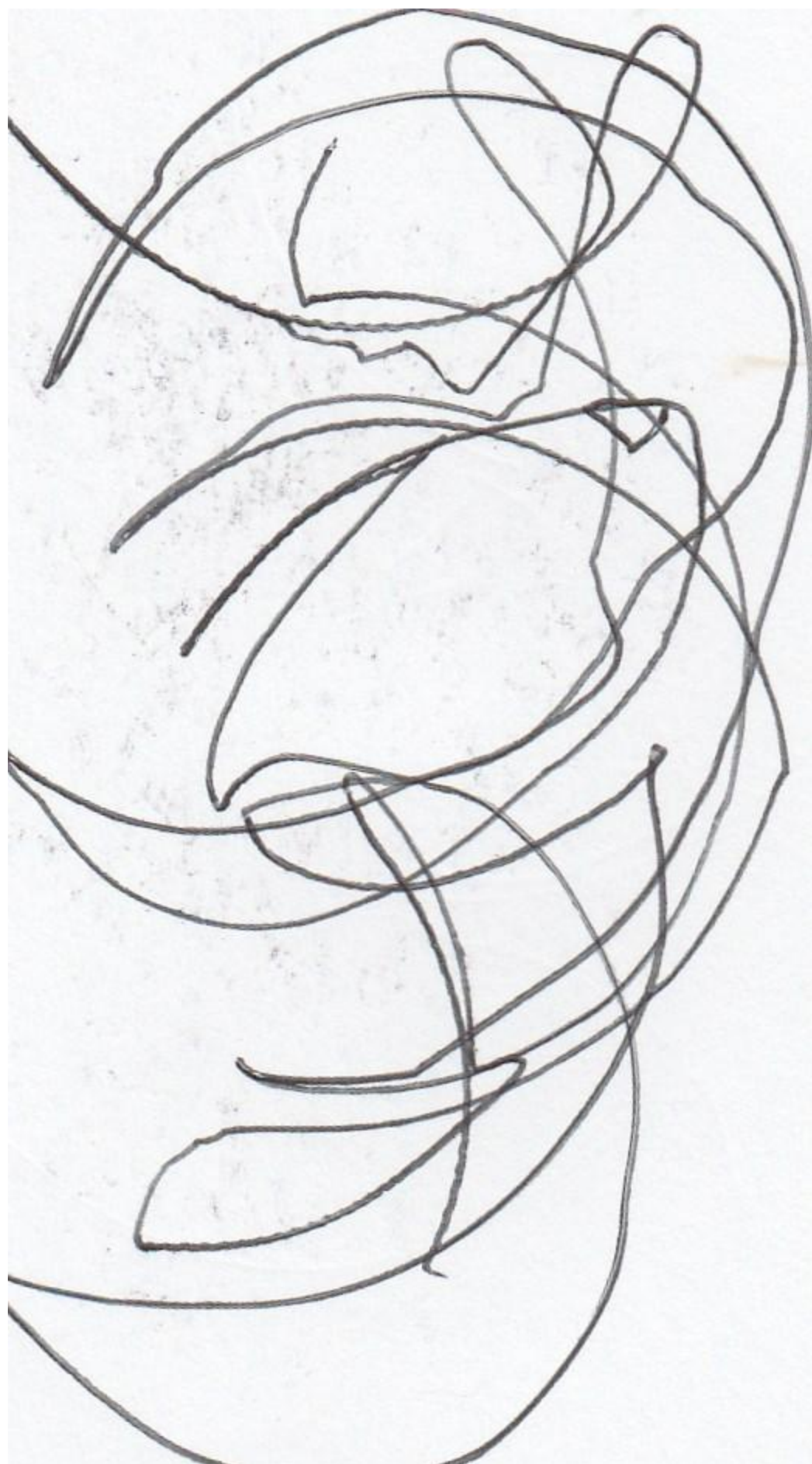
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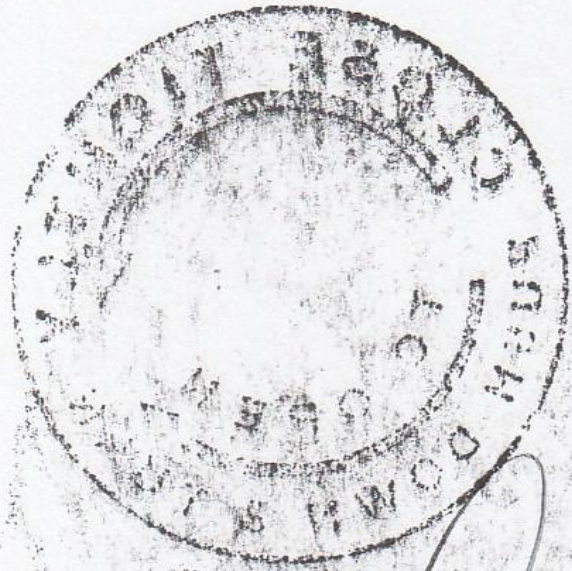
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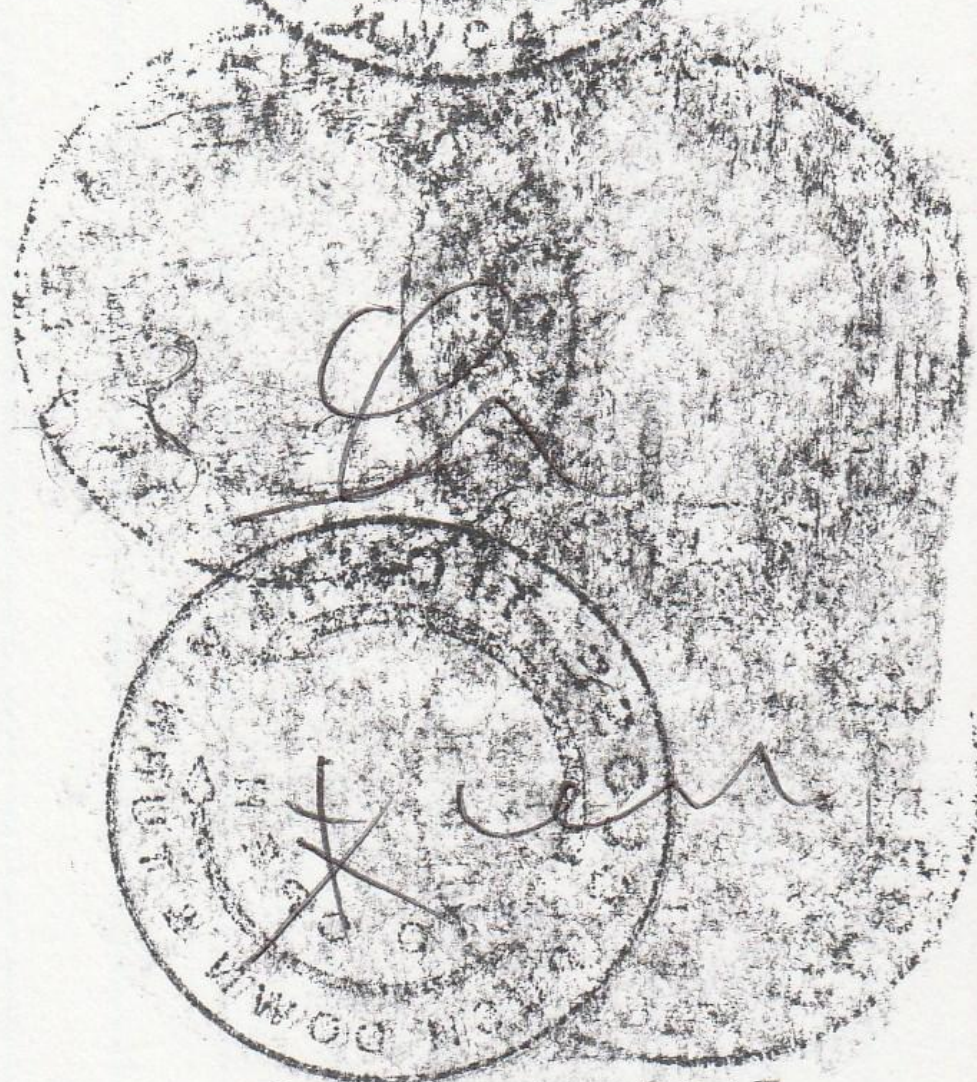
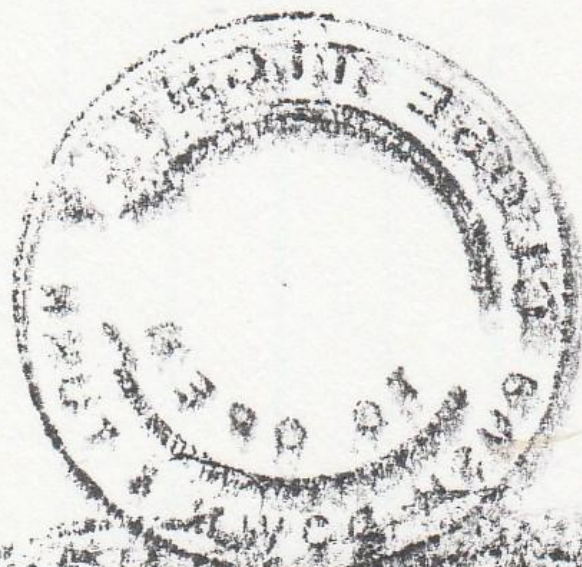
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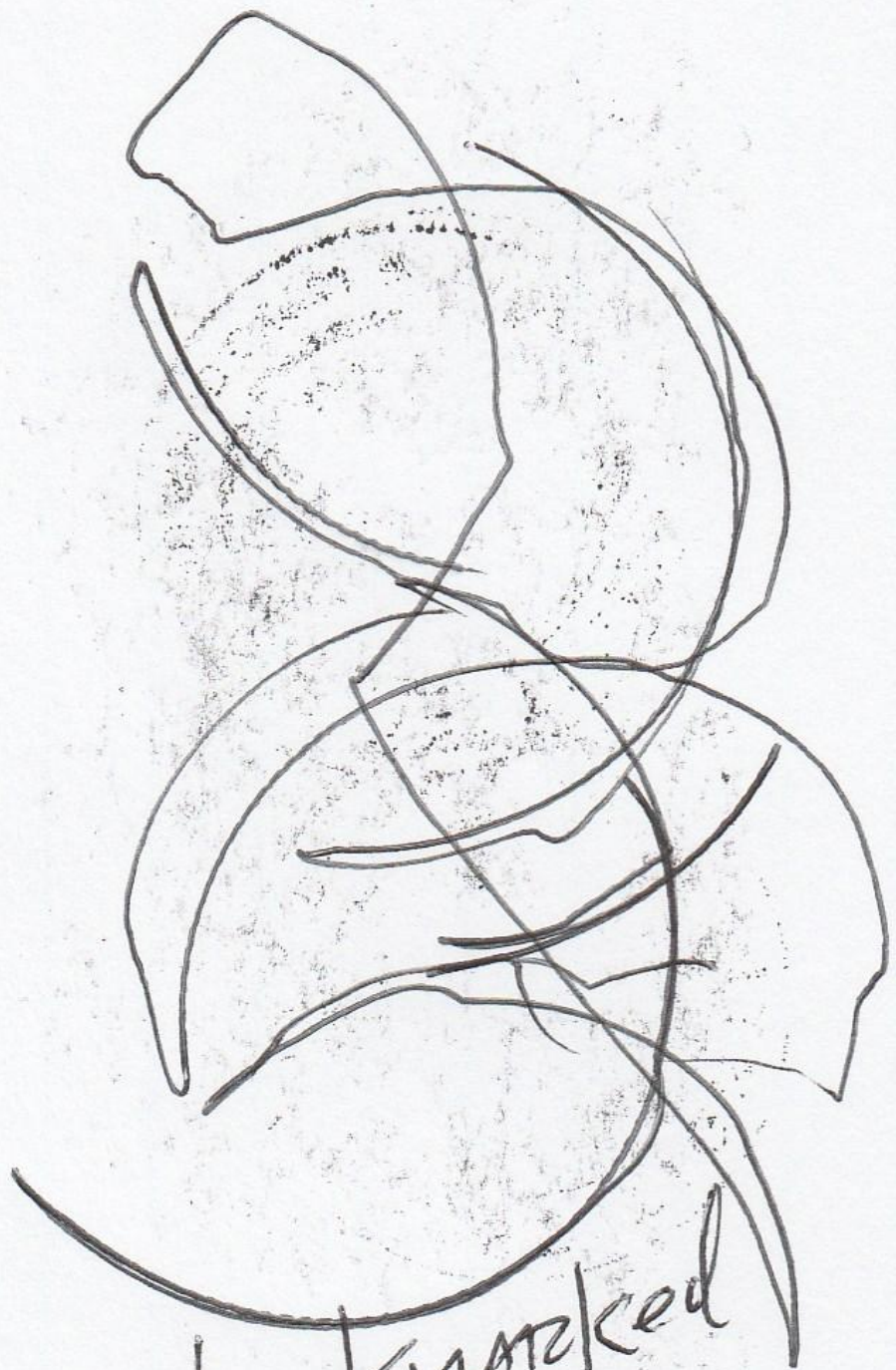
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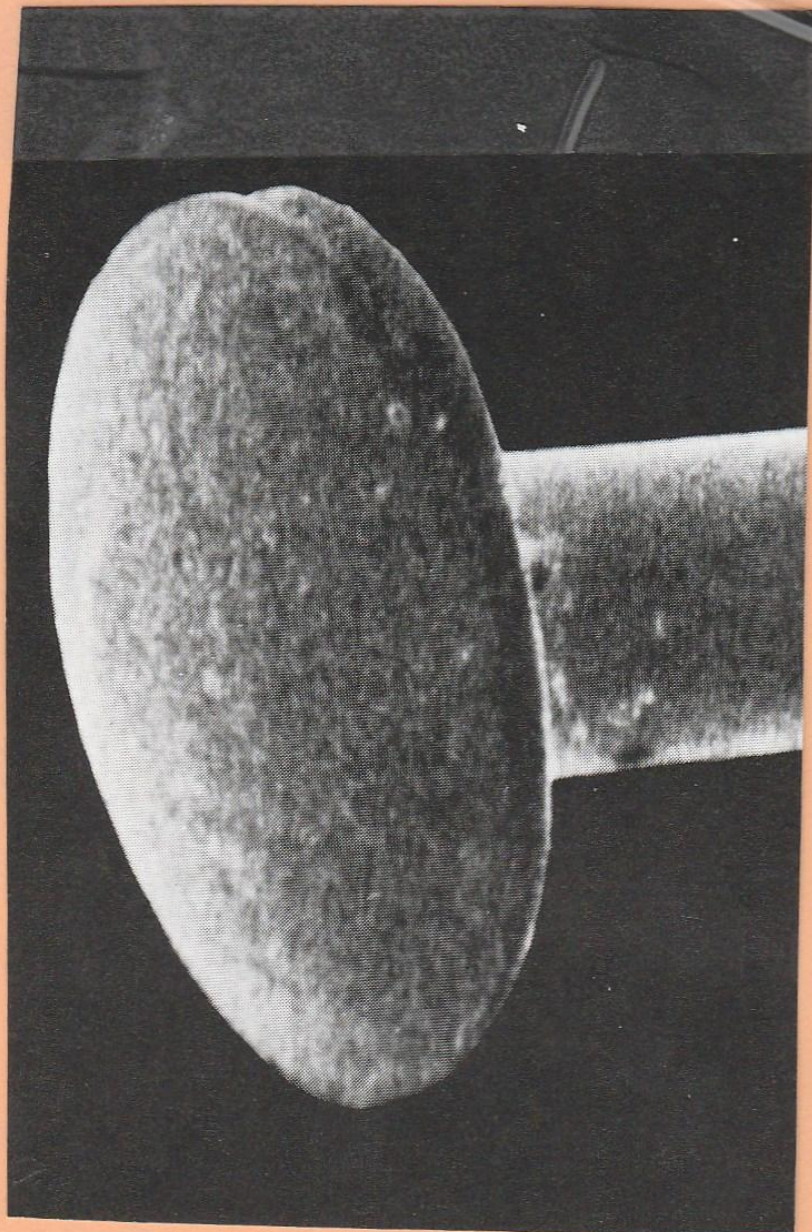
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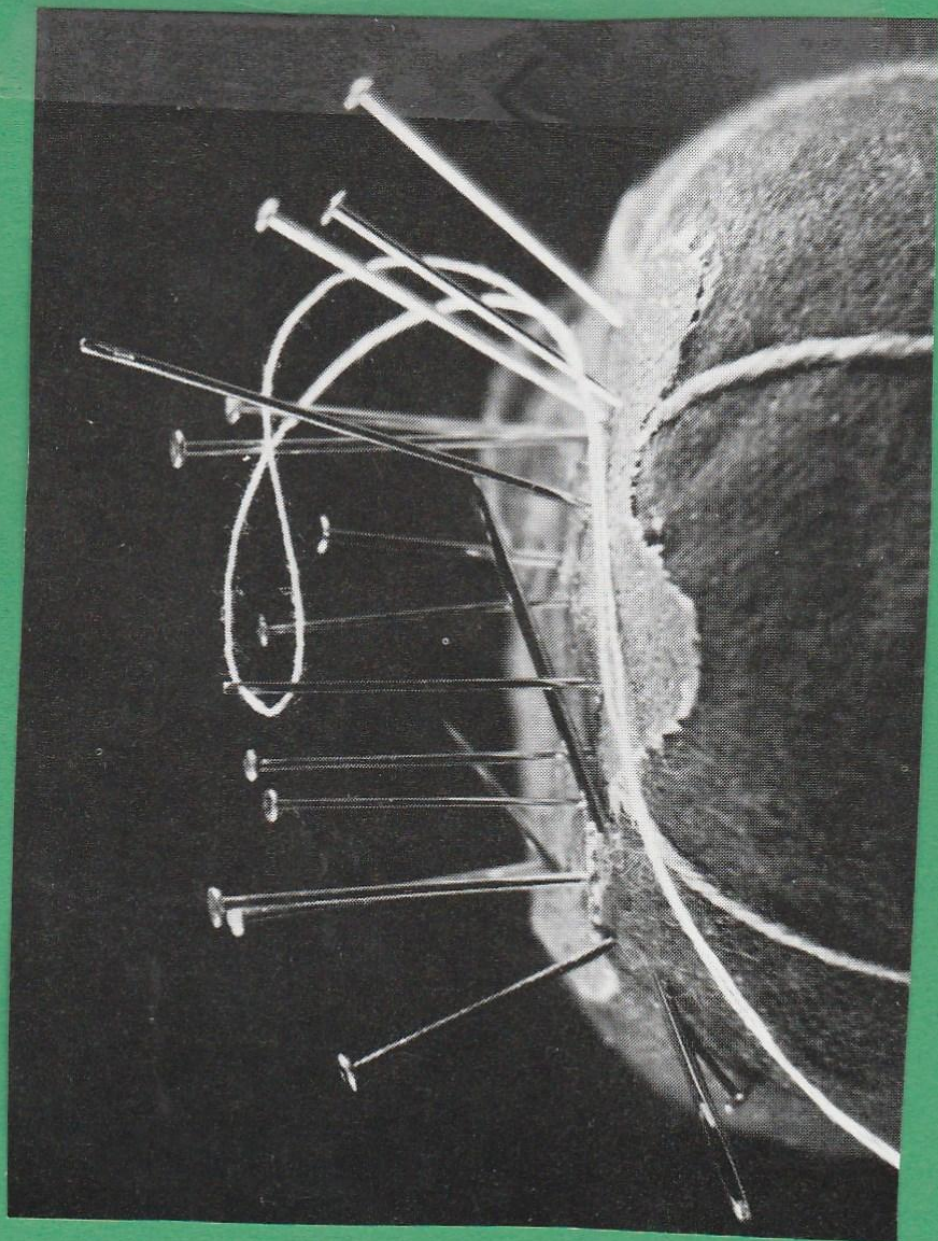
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APR 30 2017

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4411-12S	4KFH6
4411-16S	4KFH7
4411-20S	4KFH8

190265-6S	4VUP5
190265-16S	4VUP6

190261-6S	2F281
190261-12S	4KFF7
190261-16S	4KFF8

### Fittings

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yester inner braid  
steel with RoHS-

Mfr. Model	Item No.
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300-04-25	2F611
300-05-25	4VT19
300-06-25	2F613
300-08-25	2F615
300-10-25	4VTN1
300-12-25	4VTN2
300-16-25	4VTN3
300-20-25	4VTN4

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28	3	FC300-10-25
28	3 3/8	FC300-12-25
28	4	FC300-16-25
28	4 3/8	FC300-20-25
28	5 1/2	FC300-24-25
28	6 1/2	FC300-28-25
28	7 3/8	FC300-32-25
28	9	FC300-36-25
28	6	FC300-40-25

190297-00	
190297-10S	4KFF4

190296-4S	2F285
190296-5S	4KFF9
190296-8S	2F287
190296-10S	2F283

190295-4S	4KFG6
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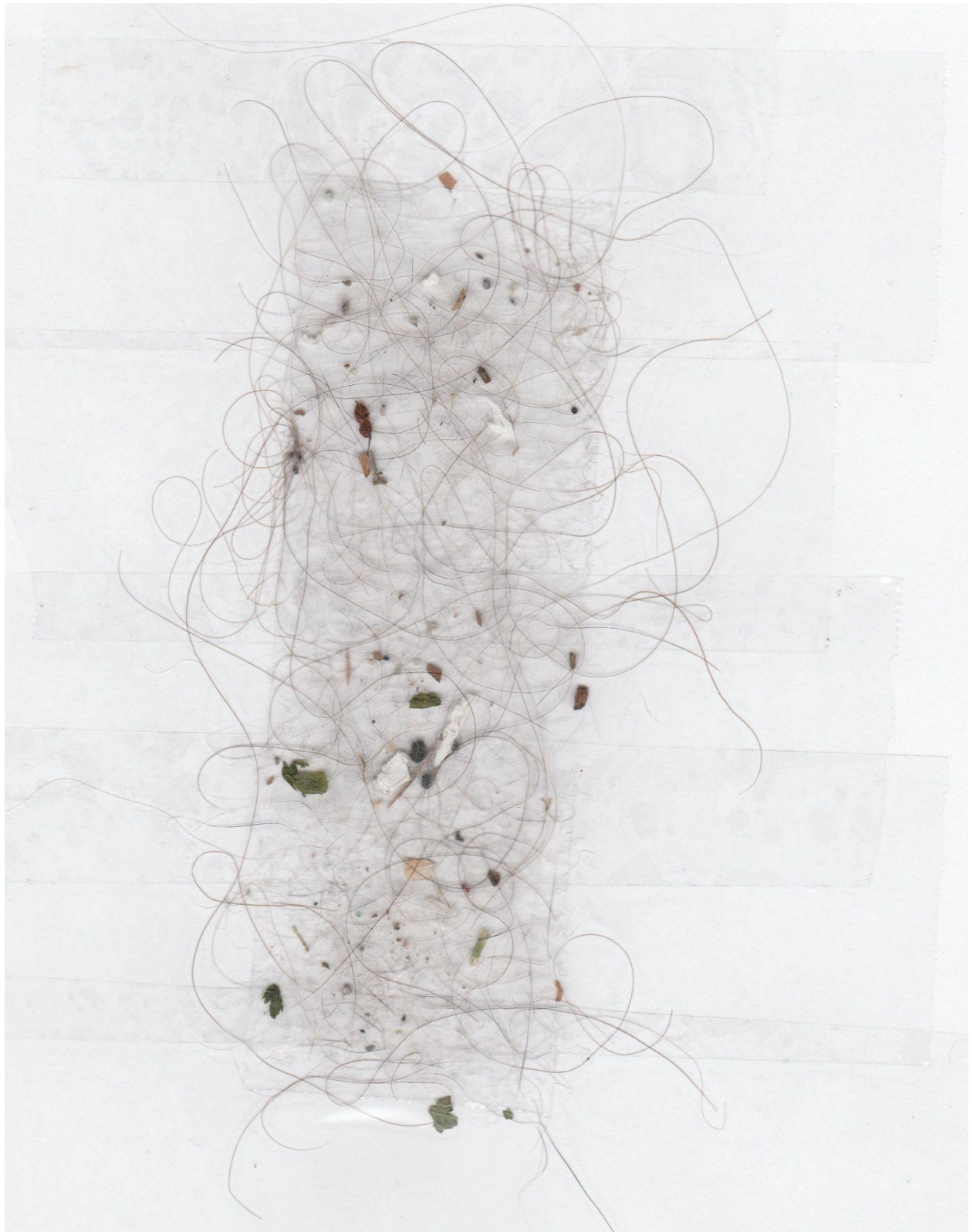
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04.26.2017

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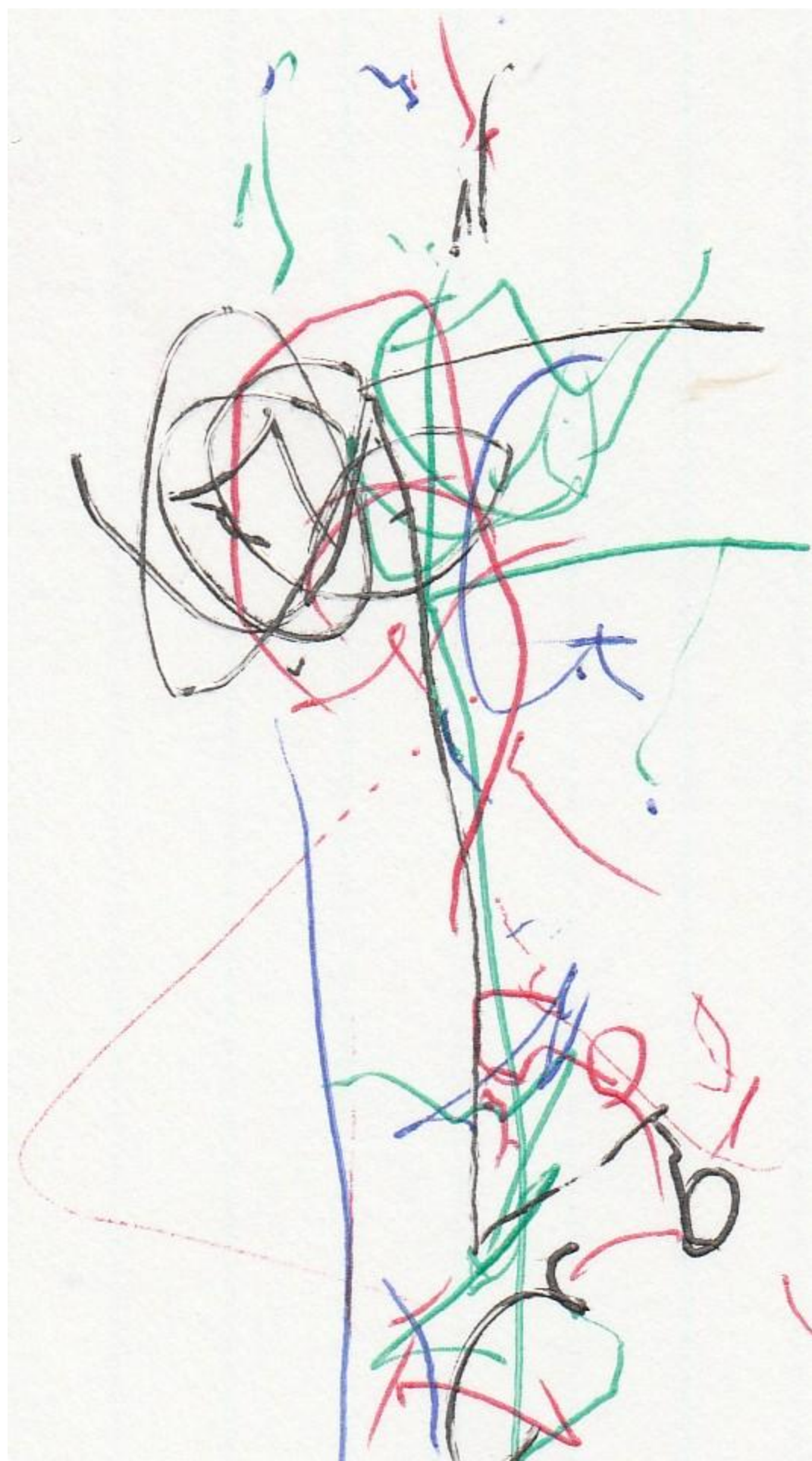
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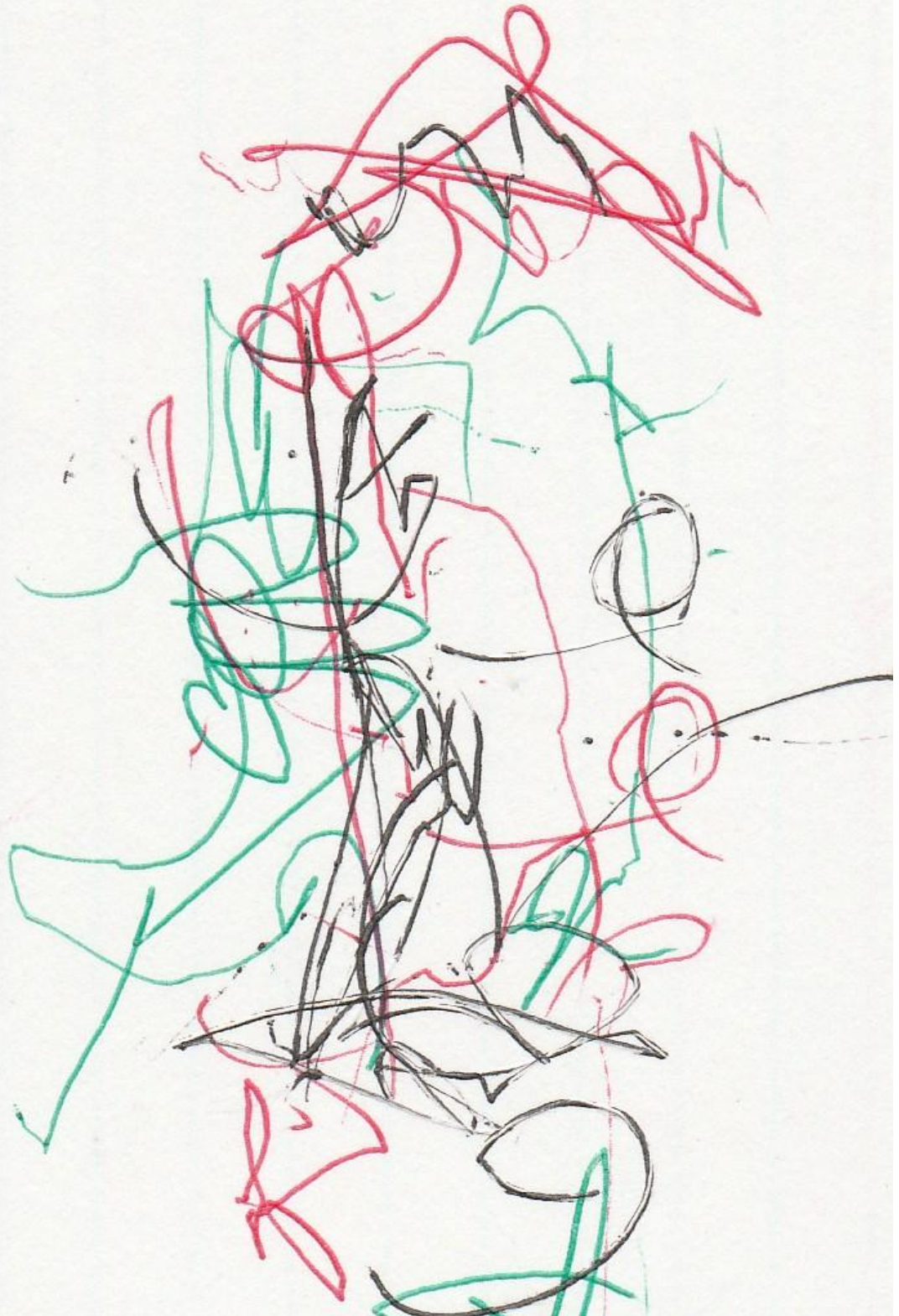
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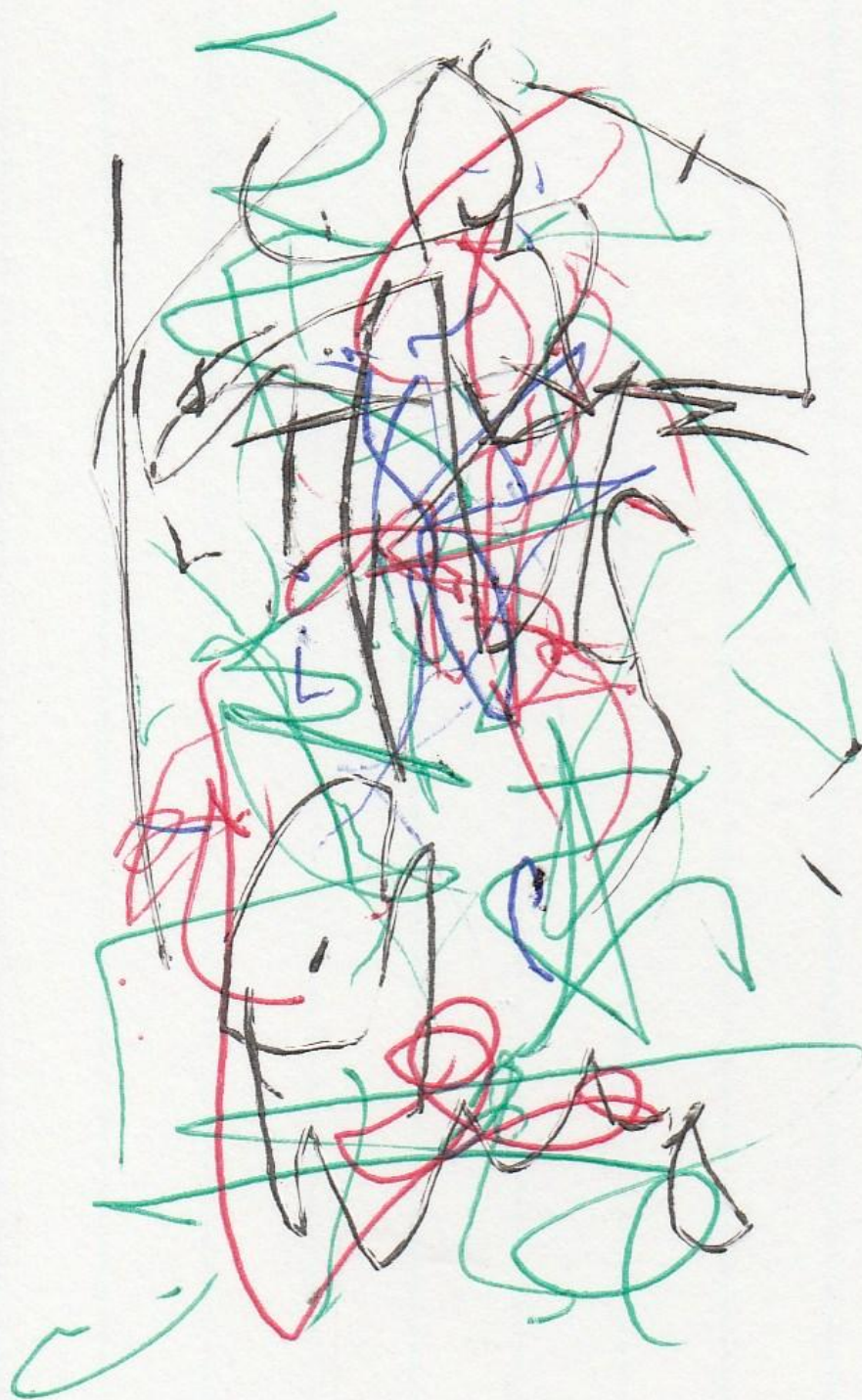




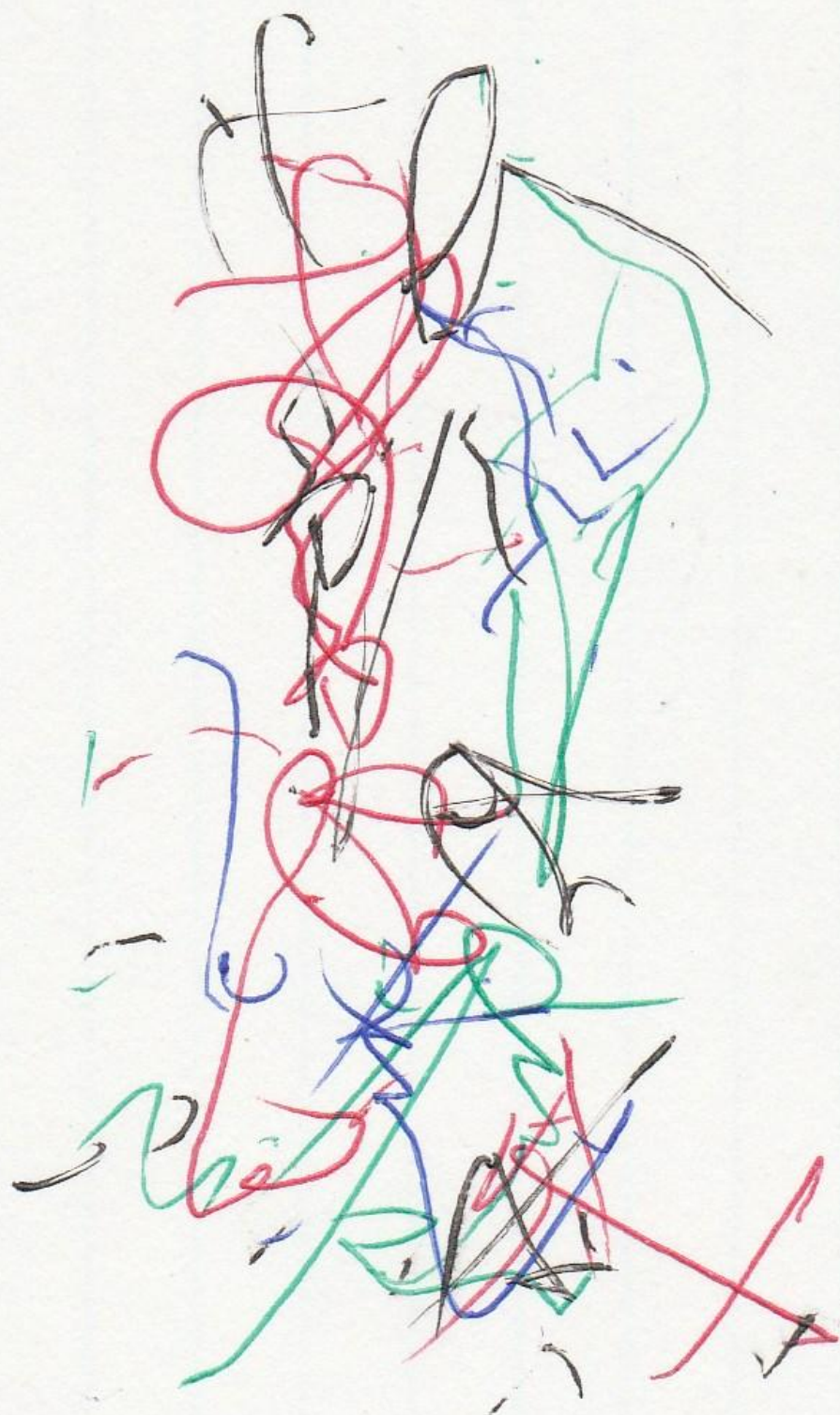
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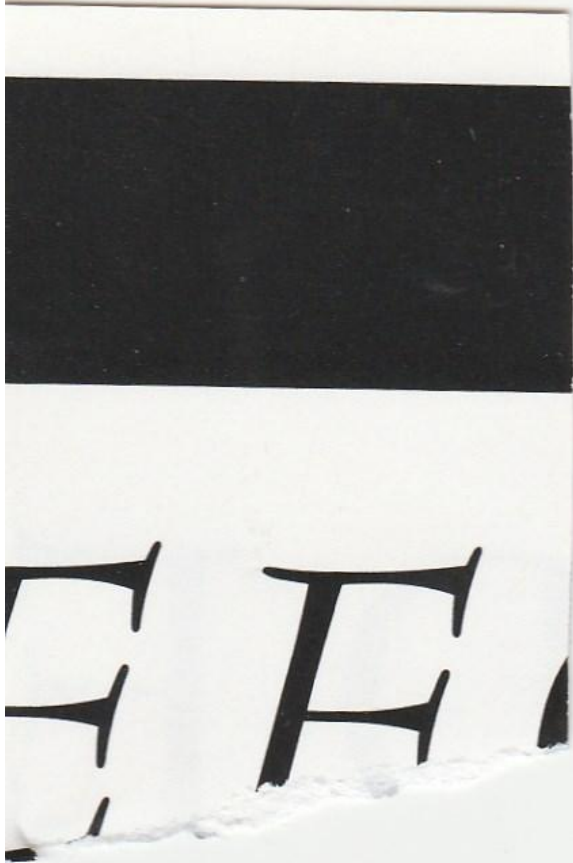


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**concrete**



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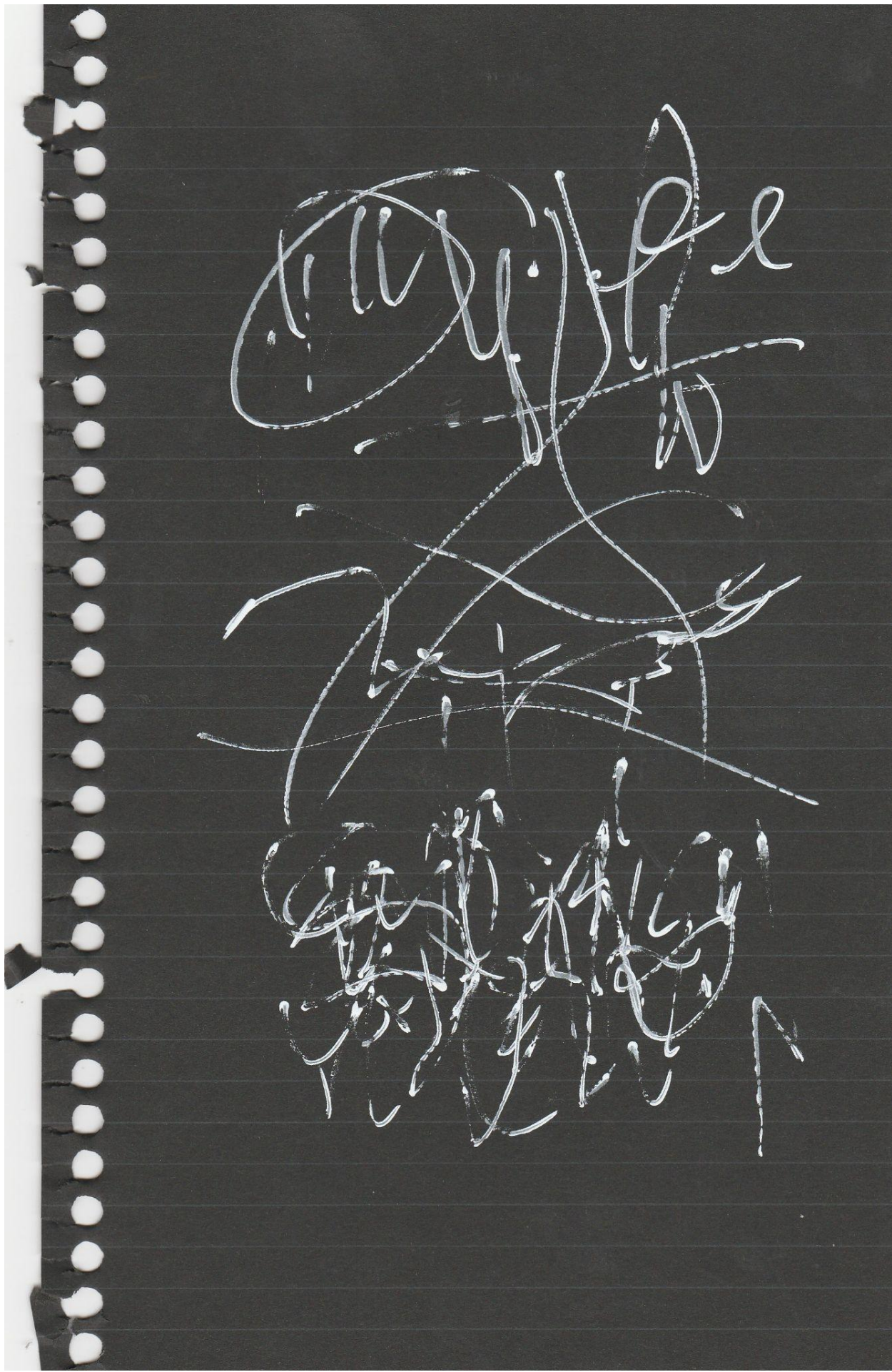


**CONTAINERS**











of course then the so easy

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of course then the so easy



of course

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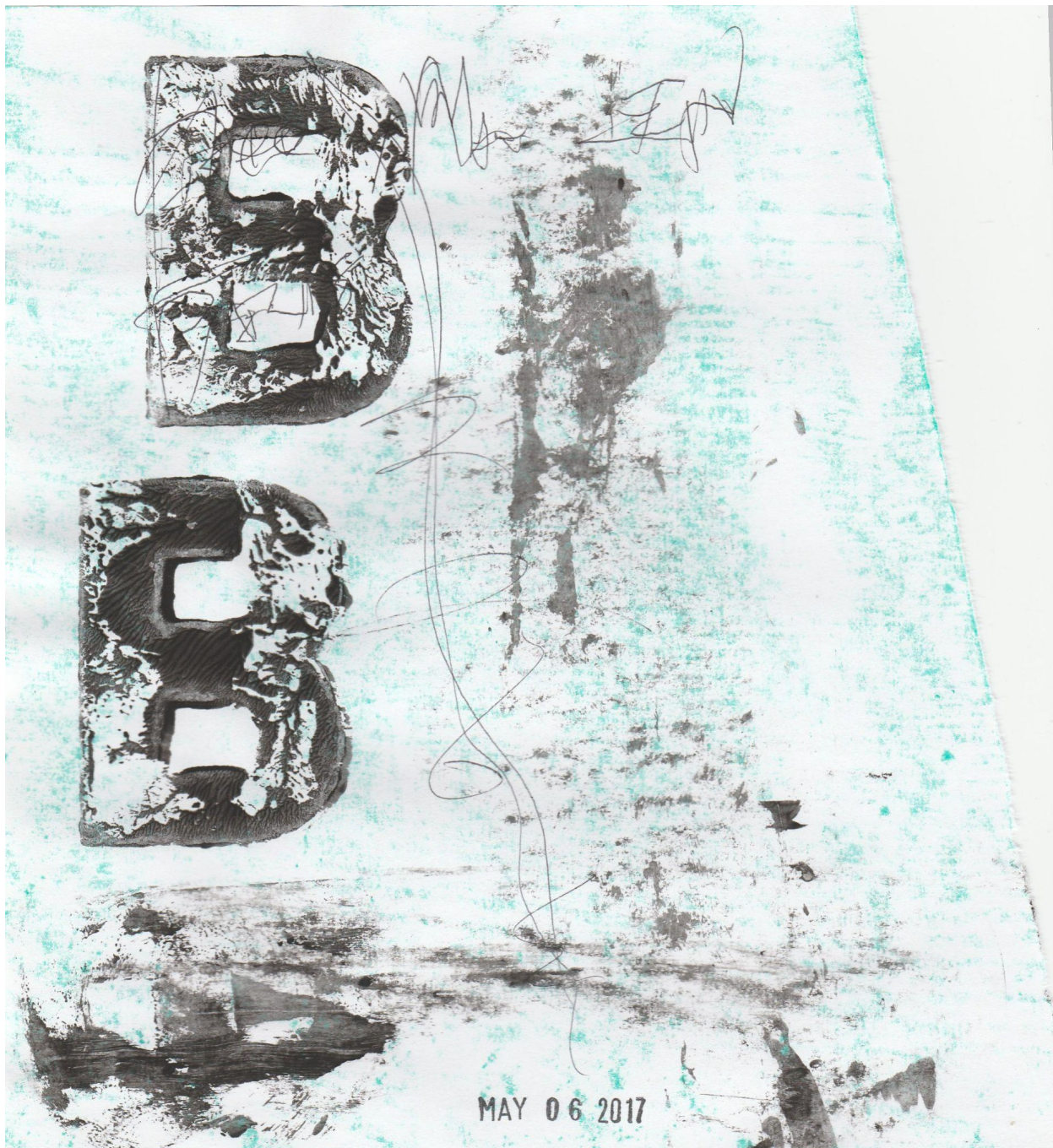
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MAY 06 2017







What it has become —



MAY 06 2017



$$\cos \beta(x - vt) = \cos(\beta vt - \beta x),$$

then  $\omega = \beta v$  is the (angular) frequency of the wave and  $\lambda = 2\pi/\beta$  is the wavelength. Furthermore,  $\beta = \omega(LC)^{1/2}$ . The typical forward wave can also be expressed as an exponential function so that

$$I_1(x,t) = ReI_1 e^{j(\omega t - \beta x)} \\ V_1(x,t) = ReV_1 e^{j(\omega t - \beta x)} \\ V_1 = Z_0 I_1$$

The reverse wave is

$$I_2(x,t) = ReI_2 e^{j(\omega t + \beta x)} \\ V_2(x,t) = ReV_2 e^{j(\omega t + \beta x)} \\ V_2 = -Z_0 I_2$$

with  
in which

If in this example the origin of the  $x$ -axis is chosen so that it is negative along the direction of the electromotive source, then

$$I = \frac{V - V_2}{Z_0} = \frac{V - (-Z_0 I_2)}{Z_0} \\ I = \frac{V + Z_0 I_2}{Z_0}$$

and the voltage

$$V = Z_0 I + V_2 = Z_0 I + (-Z_0 I_2)$$

The square of the voltage amplitude is found by multiplying the voltage ( $V$ ) and its complex conjugate ( $V^*$ ) together,

$$V V^* = (Z_0 I + V_2)(Z_0 I^* + V_2^*) = Z_0^2 I I^* + Z_0 I V_2^* + Z_0 I^* V_2 + V_2 V_2^*$$

and this has a minimum value ( $V = 0$ ) when  $V_2 = -Z_0 I$  and a maximum value ( $V = 2Z_0 I$ ) when  $V_2 = Z_0 I$ . The voltage standing ratio (vswr) is defined as the maximum to minimum amplitudes—

$$vswr = \frac{V_{max}}{V_{min}} = \frac{2Z_0 I}{0} = \infty$$

(Definitions of vswr differ; many books define it as minimum to maximum rather than maximum to minimum.) A measure of the vswr relates to the fact that this is an important technique in radio engineering.

The power ( $P$ ) dissipated in the load is  $Z_0 I^2 = (1/Z_0)(V^2 - V_2^2)$  using root-mean-square amplitudes. This can be expressed in terms of the maximum amplitude:

$$P_{max} = \frac{V_{max}^2}{2Z_0} = \frac{V_{max}^2}{2Z_0} \left( \frac{1}{vswr} \right)^2 = \frac{V_{max}^2}{2Z_0} \left( \frac{1}{vswr} \right)^2$$

$Z_0 P (vswr)^2$ . Thus, if the load is badly matched to the line and the vswr is high, it is implicitly assumed in writing [62] that  $R > Z_0$  and  $Z_0$  and  $V_{max}$  have the same sign, then  $V_{max}$  is much greater than  $V$  and there have been too many the same power transfer to a well-matched load. Because at high radio and microwave frequencies the power-handling capacity of a line is often limited by a voltage breakdown, this is a strong reason for attempting to match the load to the line.

Similar calculations can be used to obtain the voltage  $V$  at points on the line a distance  $x$  back from the load. This is, of course, the impedance that would be measured at this point. The calculation, though straightforward and tedious, and the result is

$$Z = \frac{V}{I} = Z_0 \frac{R \cos \beta l + j \sin \beta l}{Z_0 \cos \beta l + j R \sin \beta l}$$

Here,  $\beta l = 2\pi/\lambda$  so that  $\beta l/\pi$  is the length of the line measured in wavelengths.

Whenever  $\beta l = n\pi$  ( $n = 0, 1, 2, 3$ , etc.)—i.e. the line is an integral number of half wavelengths long—this gives  $\sin \beta l = 0$  and  $Z = R$  again; but if  $\beta l = (n + 1/2)\pi$ —i.e. the line is an odd number of quarter wavelengths long—this gives  $\cos \beta l = 0$  and  $Z = Z_0^2/R$ . The larger  $R$  is made, the lower is the value of  $Z$  measured at the other end of the line. A quarter wavelength of line acts as an impedance transformer. A cable of characteristic impedance  $Z_0$  and a load  $R$  can be matched by inserting a quarter wavelength of line of characteristic impedance  $Z_0'$  so that  $Z_0 = Z_0' = \sqrt{Z_0 R}$ . In other words,  $Z_0'$  is the geometric mean of  $Z_0$  and  $R$ . This is a useful technique in microwave

electronics. It may also be noted that it is essentially the principle underlying the use of thin surface layers (lens coating, or blooming) to reduce reflection of light waves from optical components, such as camera lenses. In the case when  $R$  is equal to zero—i.e. the line is shorted at the end—then

$$Z = jZ_0 \tan \beta l \quad (63)$$

When  $\beta l$  is small,  $\tan \beta l = \beta l$  and this gives

$$Z = j\omega L = j(L/C)^{1/2} \omega (LC)^{1/2} l = j\omega Ll$$

In other words, the line appears to have an inductance  $Ll$ , which is hardly unexpected because  $L$  is the inductance per unit length. If  $R$  is infinite—i.e. the line is open circuited—the results

$$Z = \frac{Z_0}{j \tan \beta l}$$

and, if  $\beta l$  is small,

$$Z = \frac{Z_0}{j \beta l} = \frac{1}{j\omega Cl}$$

The line behaves in this case like a capacitance  $Cl$ . The effect resulting from the use of cables having lengths that are a small fraction of a wavelength can be treated solely in terms of inductance  $Ll$  and capacitance  $Cl$ . This is the standard low-frequency approximation.

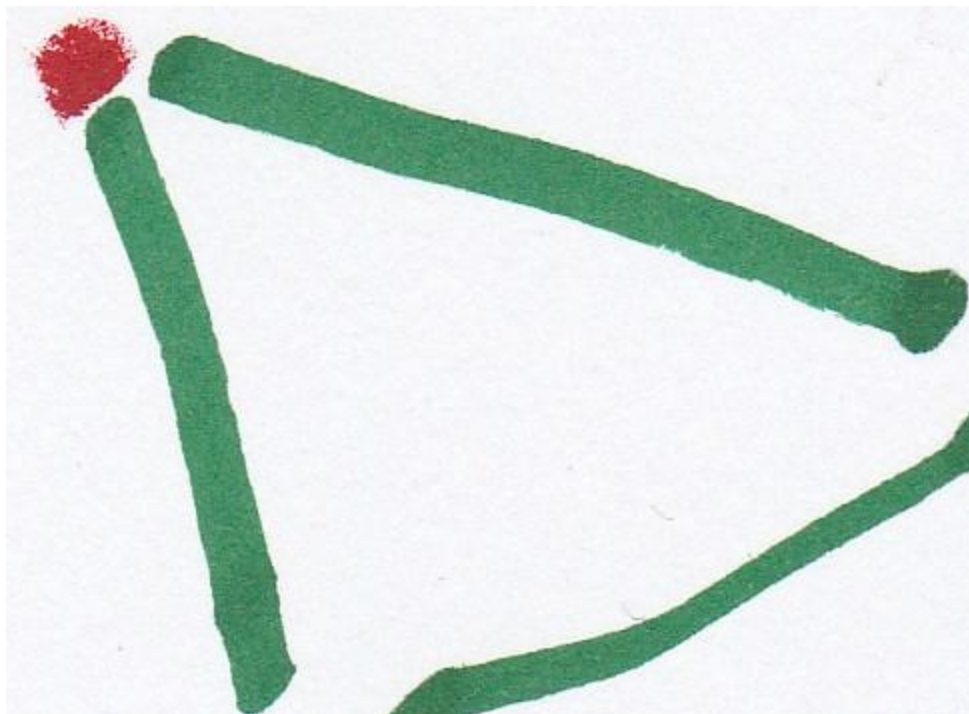
When  $\beta l = \pi/2$ , equation (63) shows that  $Z \rightarrow \infty$ . Thus a shorted line one-quarter wave long appears at the other end as an open circuit. As the frequency is changed so that  $\beta l$  increases through  $\pi/2$ , the phase angle of  $Z$  changes from  $+\pi/2$  to  $-\pi/2$  and  $|Z|$  goes through a sharp minimum. This is much like the behaviour of a tuned circuit. In microwave electronics, short-circuited lines are used as substitutes for tuned circuits. Although lines have resistive loss and this broadens the resonance, the effective quality factor  $Q$  is still high. For a 1000 ohm resistor, when the wavelength is 10 centimetres (4 inches), the line (one inch) long behaves as a resonant circuit. This example demonstrates the way in which the transmission delay in cables makes microwave circuit analysis quite different from low-frequency ac circuit analysis.

Real transmission lines always have a finite resistance  $R$  per unit length. If finite resistance is included in the theory, the equation that replaces (56) is

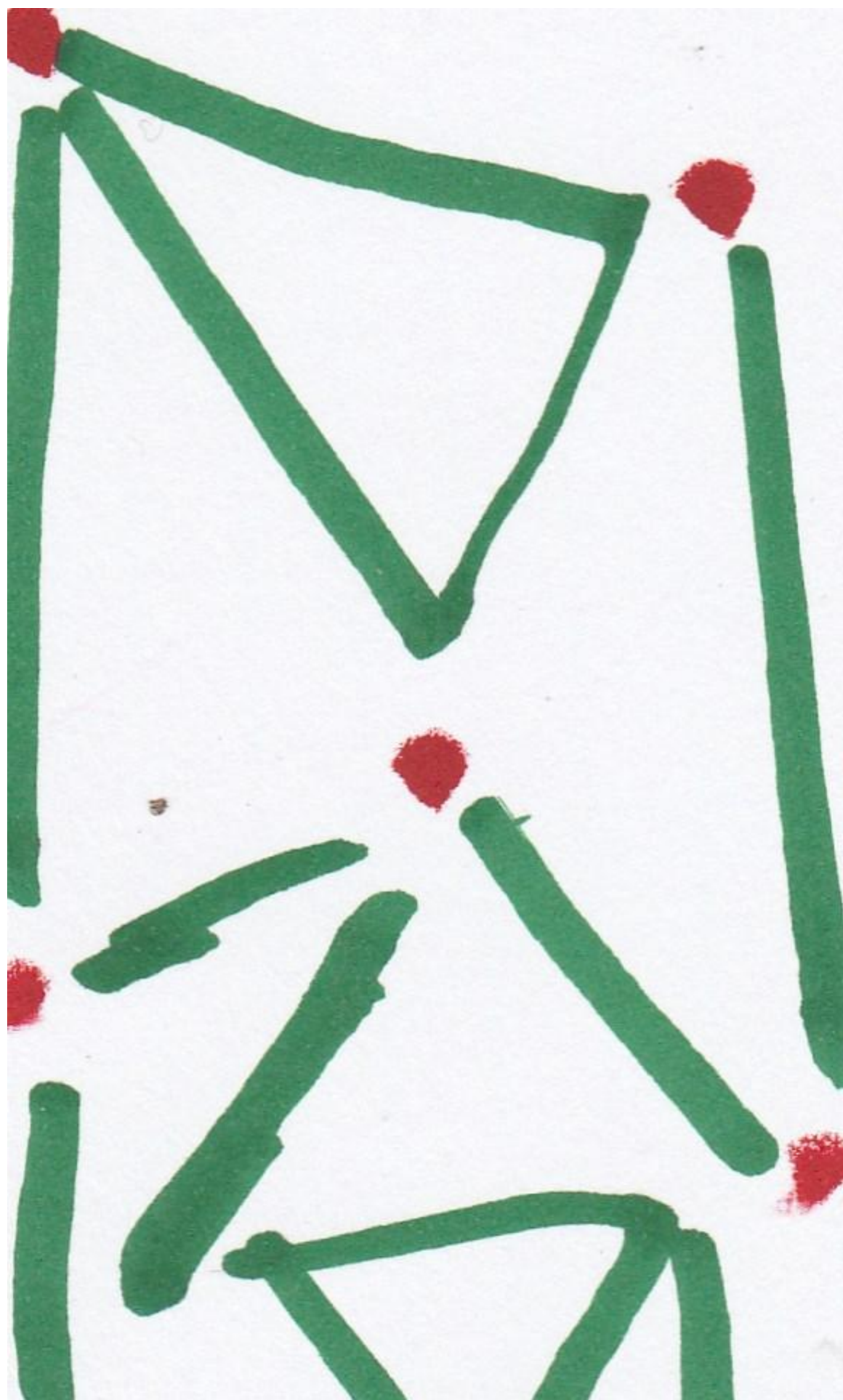
$$\frac{\partial V}{\partial x} = -\frac{\partial I}{\partial x} = -RI$$

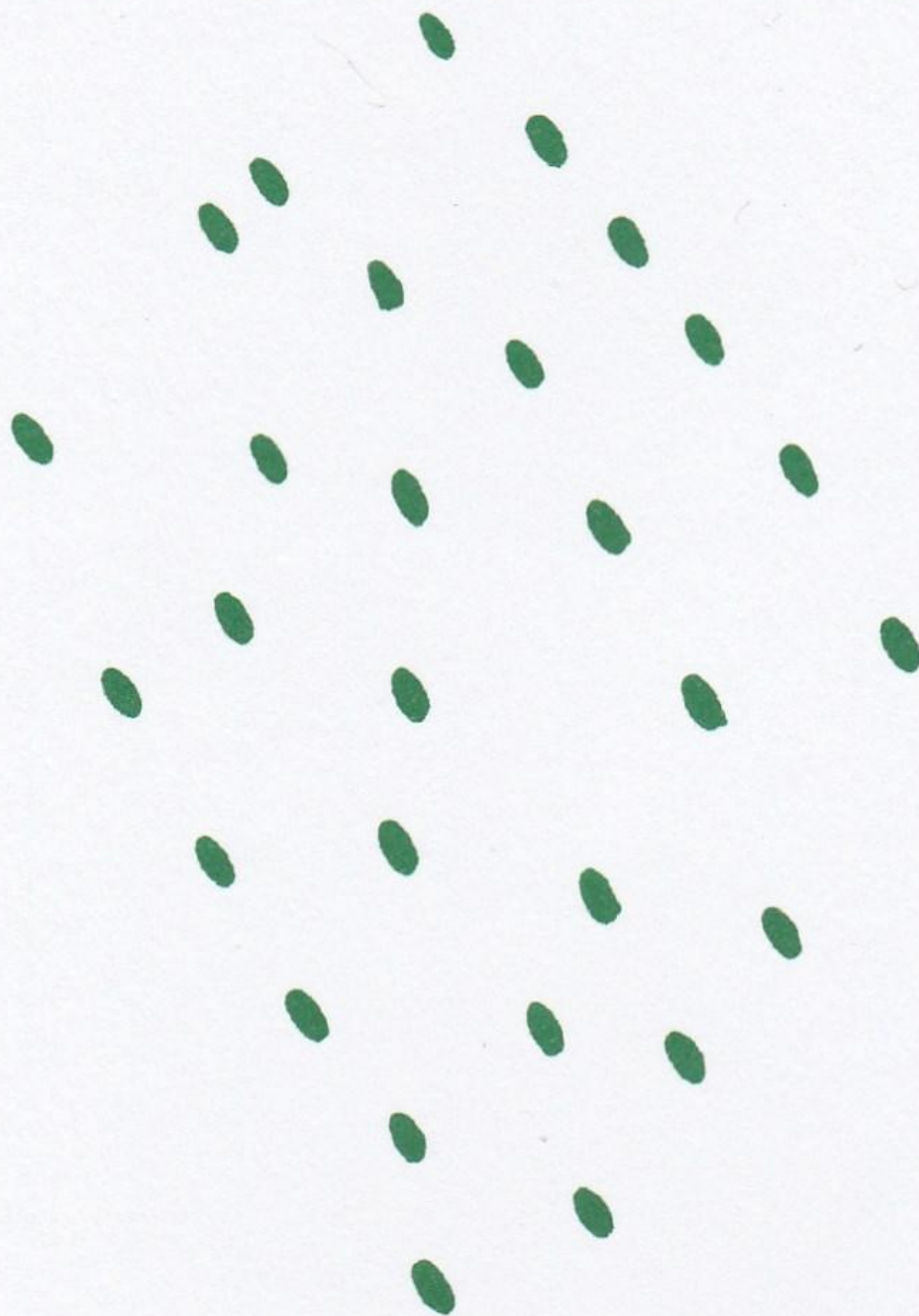
and the wave, instead of propagating unattenuated as  $e^{j(\omega t - \beta x)}$ , is  $e^{j(\omega t - \beta x) - ax}$  in which  $a = R/2Z_0$ . The value of  $R$  is determined by the so-called skin effect (see below) and increases with frequency as  $\omega^{1/2}$ . At 100 megacycles, a frequency used for radio and television transmissions,  $a$  is of the order of  $3 \times 10^{-3}$  metre<sup>-1</sup>. The amplitude of the voltage and current on the line decrease exponentially along the line as  $e^{-ax}$  and the power as  $e^{-2ax}$ . In the laboratory, with a few metres of line, the decrease is not so great. For example  $x = 10^4$  metres (six miles), the power is reduced by a ratio  $e^{-60} \approx 10^{-26}$ . If  $10^6$  watts are fed into one end of the line, a power of only  $10^{-20}$  watts emerges at the other end, and this would be barely detectable. Attenuation by the skin effect in cables at radio frequencies, although at first sight small, is a serious bar to their use in distributing radio signals over extended paths. Practical cable-distribution systems rely on amplifiers to boost the signal. These amplifiers have to be placed at frequent intervals in the cable.

It will be seen then that, at fairly low frequencies in which cables and connections are only a fraction of a wavelength long, the time delay in the cable can be more or less neglected and at most an allowance made for the total capacitance and inductance of the cable. At high frequencies, however, the time delay introduces quite new effects. The theory of the propagation of waves on cables is a highly developed branch of electrical engineering, but it has a more general significance in many



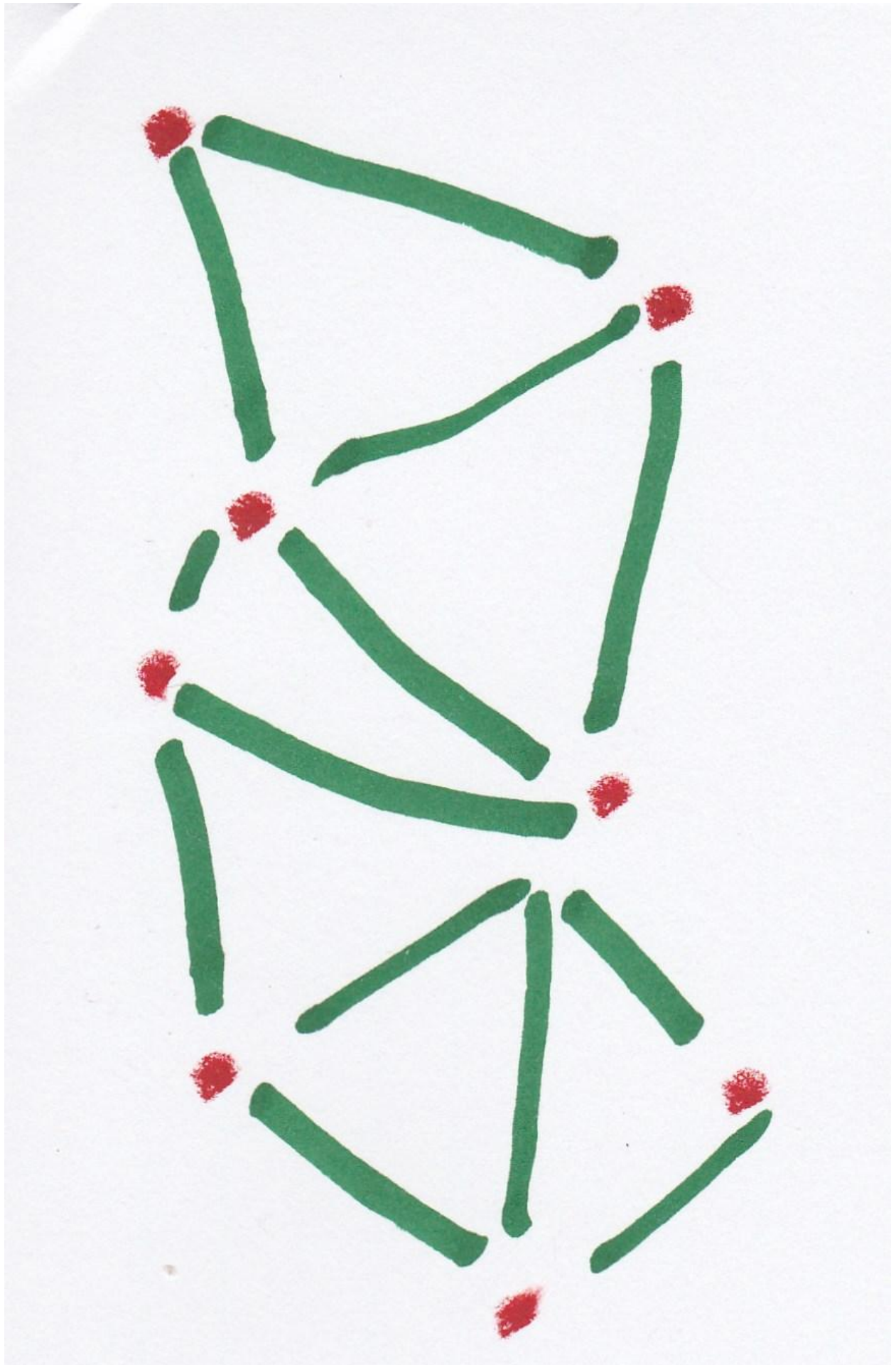






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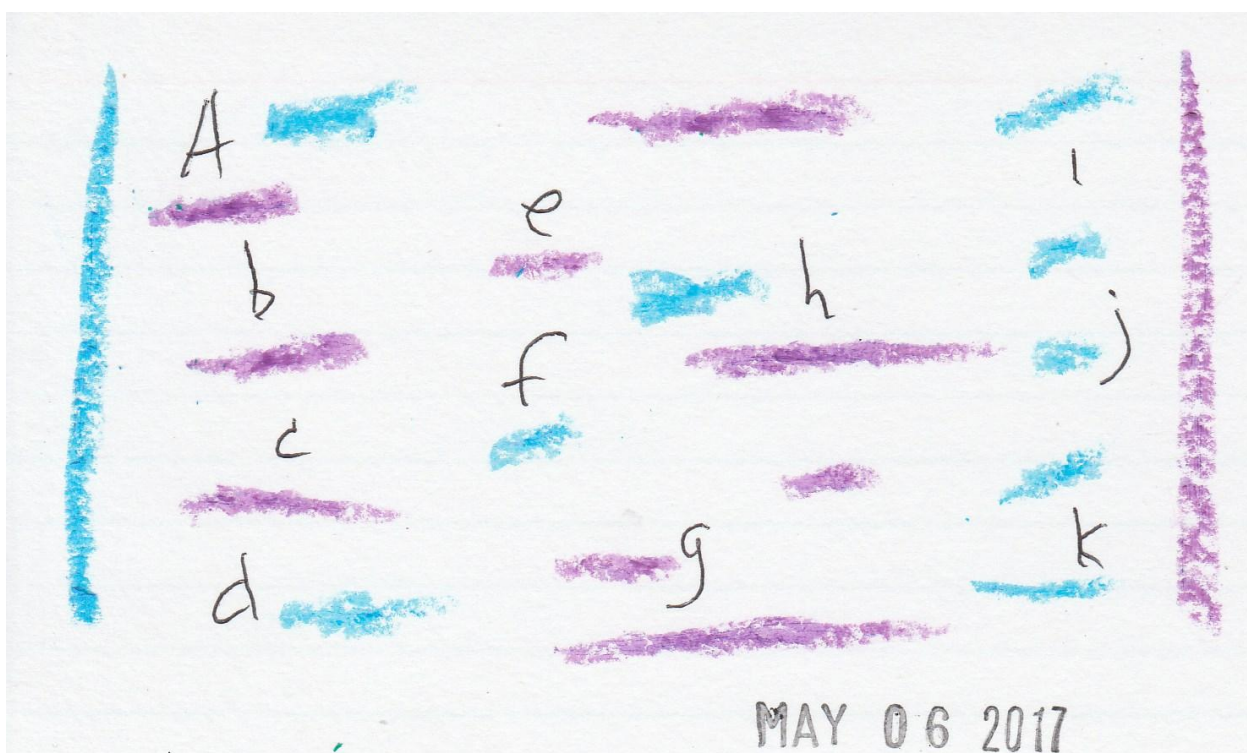
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